

Vertical-mode decomposition as a dynamic description of the Atlantic meridional overturning circulation

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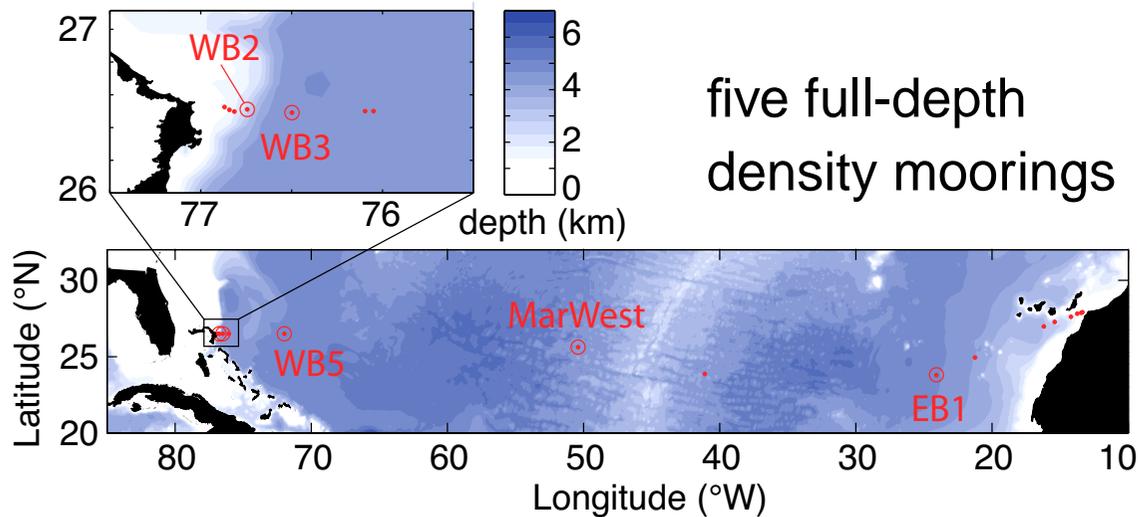
US AMOC Meeting, Miami

7 June 2010



What processes does the RAPID/MOCHA array measure?

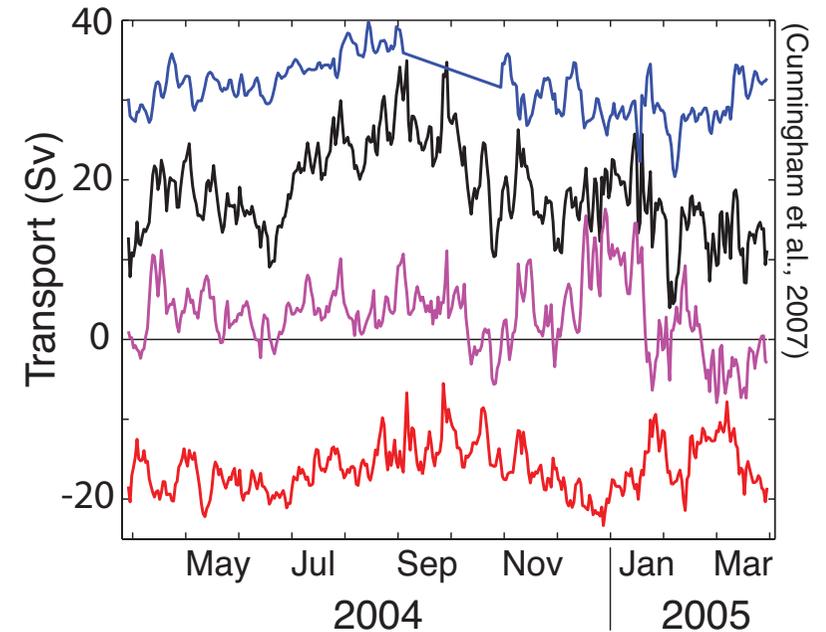
RAPID/MOCHA array



five full-depth
density moorings

AMOC Components

$$T_{AMOC} = T_{MO} + T_{FC} + T_{Ek}$$

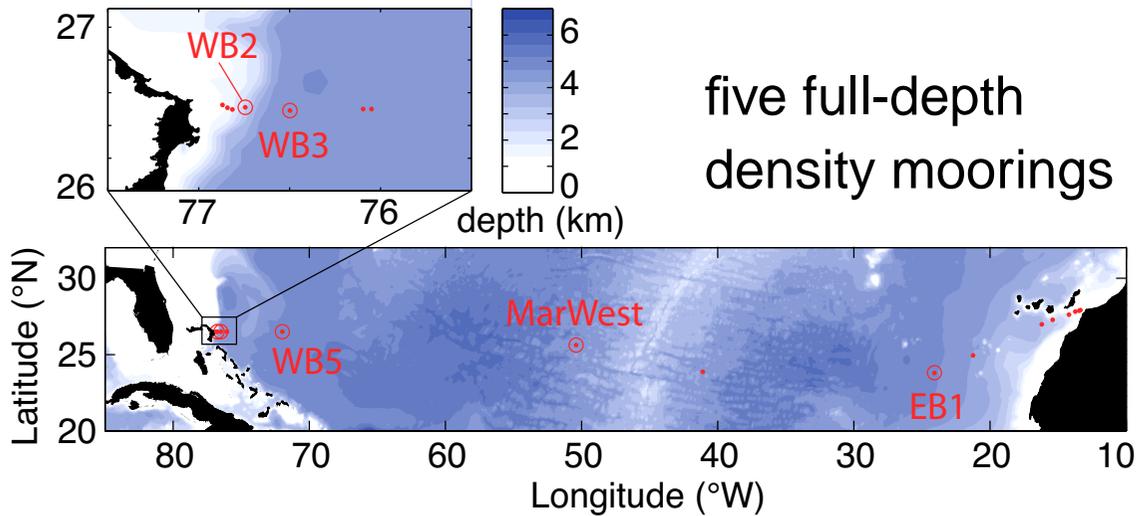


To understand the basin-scale signal (AMOC),
we must also understand the local signal.

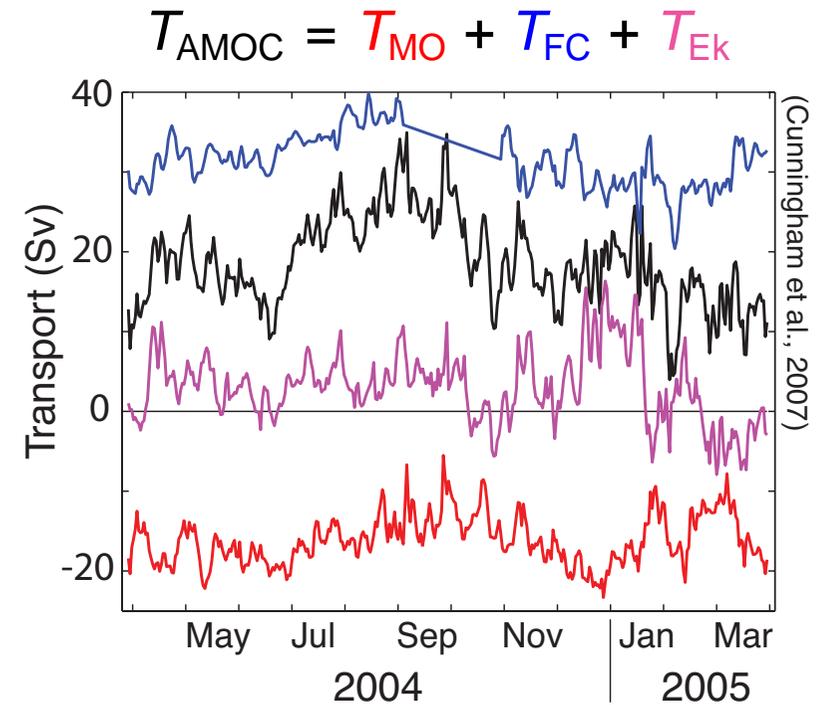
Test whether simple dynamics explain vertical density fluctuations.

What processes does the RAPID/MOCHA array measure?

RAPID/MOCHA array



AMOC Components



To understand the basin-scale signal (AMOC), we must also understand the local signal.

Test whether **simple** dynamics explain **vertical** **density** fluctuations.

use *a priori* shapes, not EOFs

vertical modes

certain forms implicitly assumed by Wunsch (2008) or for SSH

profiles of local geopotential anomaly

Methodology and Organization

1. **Decompose vertical structure at each mooring into modes.**
Fit modes to 2-day low-pass filtered moored CTD measurements:
2. **How useful is the decomposition and what does it show?**
Compare it with the original signal and with SSH.
 - local signal at each mooring
 - transports between moorings, across basin
3. Summary

How to obtain modes from our observations?

In a motionless and flat-bottomed ocean, vertical modes depend on $N^2(z)$.

$$u' = \sum_n U_n(x, y, t) F_n(z)$$

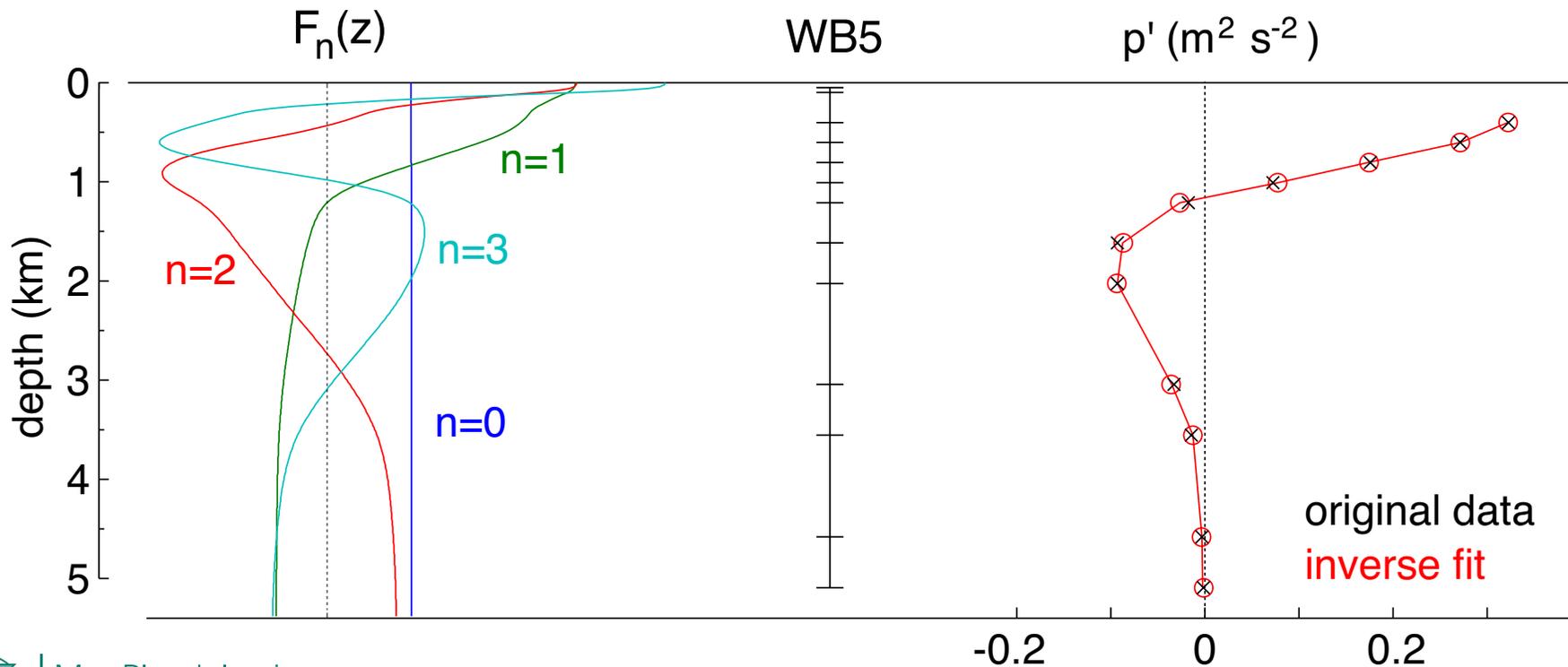
$$p' = \rho_0 \sum_n P_n(x, y, t) F_n(z)$$

From moored density measurements:

- (1) calculate $\rho(z, t)$ with climatology (RAPID)
- (2) calculate pressure perturbation

$$p'(z, t) = \frac{g}{\rho_0} \int \rho'(z, t) dz$$

- (3) use Gauss-Markov inversion to fit 8 modes at each time-step (Wunsch, 1996)



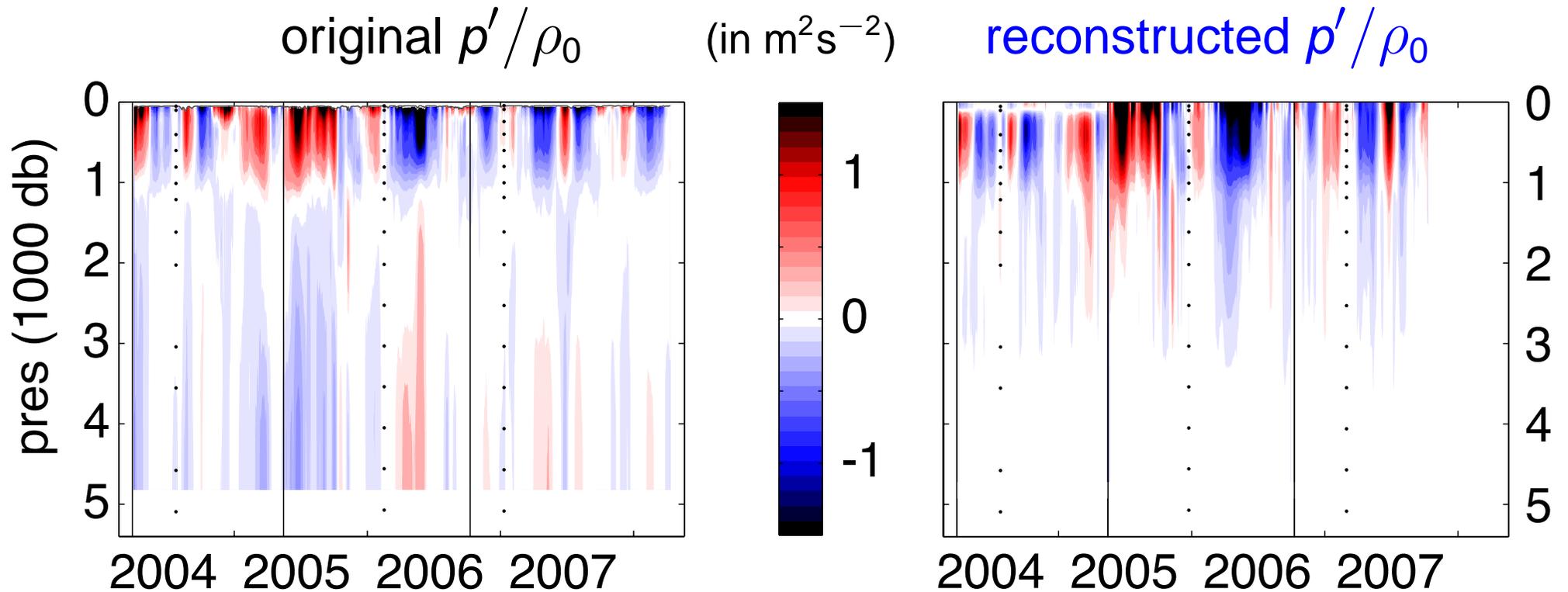
Compared quantities

For remainder of talk, compare these three quantities in units of m^2s^{-2} :

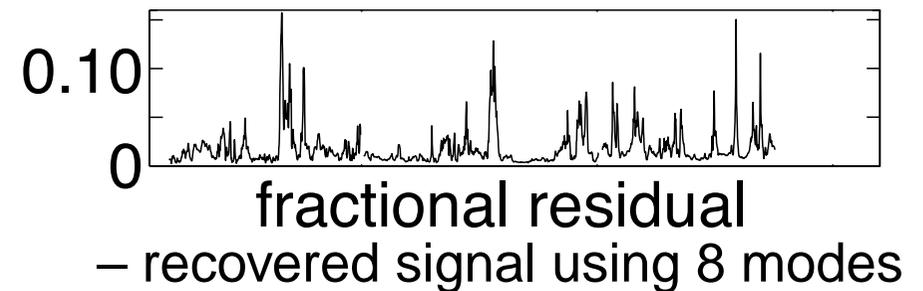
Original	geopotential anomaly perturbation:	ϕ'
	or reduced pressure perturbation:	p' / ρ_0
Reconstructed	reconstructed pressure perturbation:	$\sum_n P_n(x, t) F_n(z)$
SSH	altimetric SSH anomaly times gravity: (with zonal average removed)	$g\eta$

Because these quantities are all time-perturbations, for geostrophic transport calculations we must add back the time-averaged geopotential anomaly $\overline{\phi}(z)$.

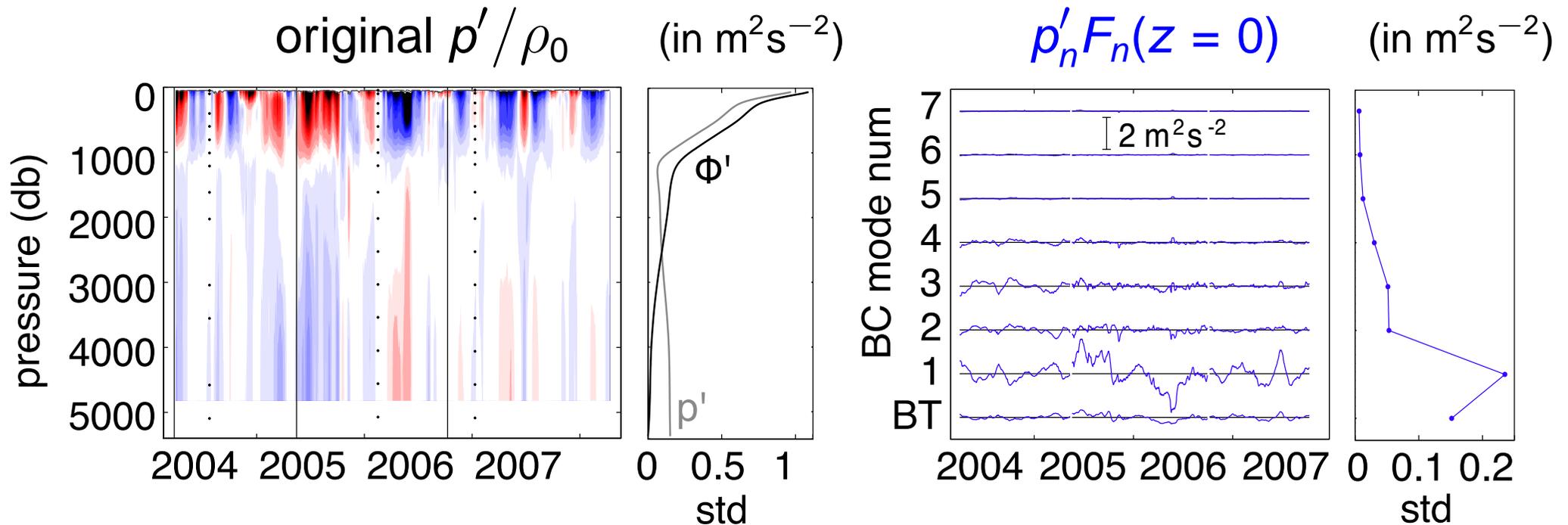
Is the full-depth signal recovered at WB5?



- data from three deployments
- only use data at sensor depths (dots)
- don't use observations above 200 m



Variance against depth or mode number

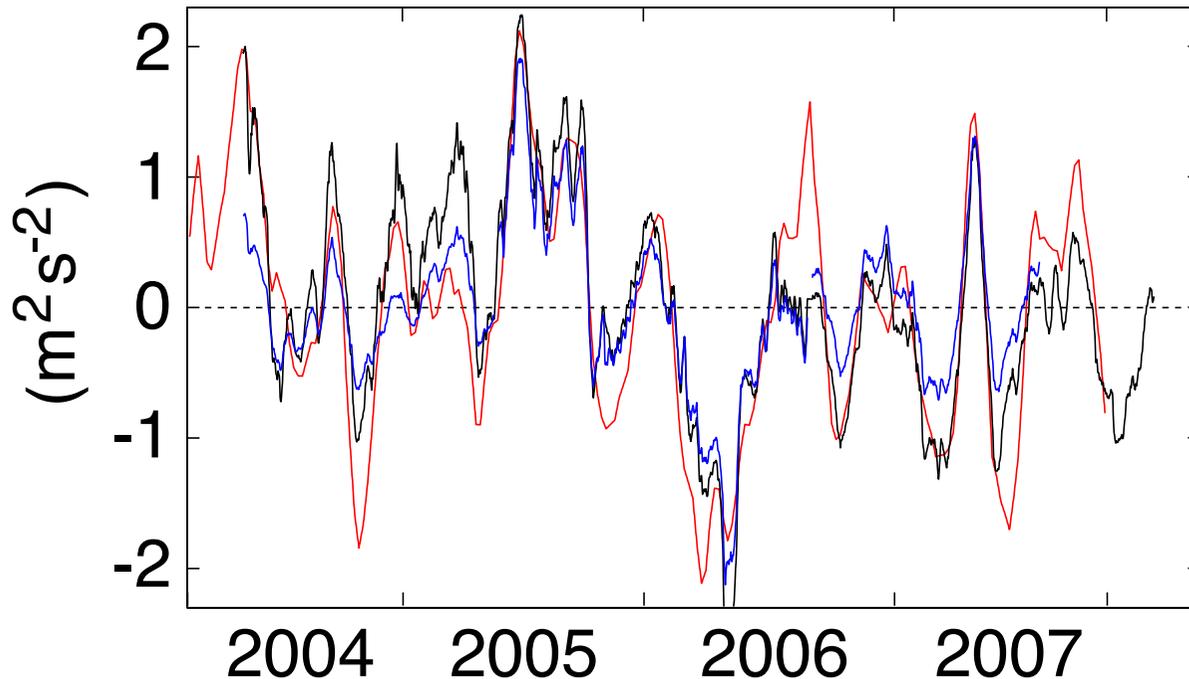


- p' is intensified above 1200 m
- weak signal below 3000 m

- first baroclinic mode (BC1) dominates
- barotropic mode (BT) is retained, as it's unclear whether enforcing a condition of 'no forcing' is appropriate (Kunze *et al.*, 2002)

Is the local signal at WB5 recovered?

geopotential anomaly (Φ) at 200 db
 BC1 reconstruction (ρ'_1/ρ_0) at 200 db
 SSH ($g\eta$) from altimetry



correlation (r)

	Φ	ρ'_1/ρ_0	$g\eta$
Φ	1	0.87	0.85
ρ'_1/ρ_0		1	0.93
$g\eta$			1

there's good agreement:

- BC1 recovers 76% of variance
- of which almost all (86%) is coherent with SSH

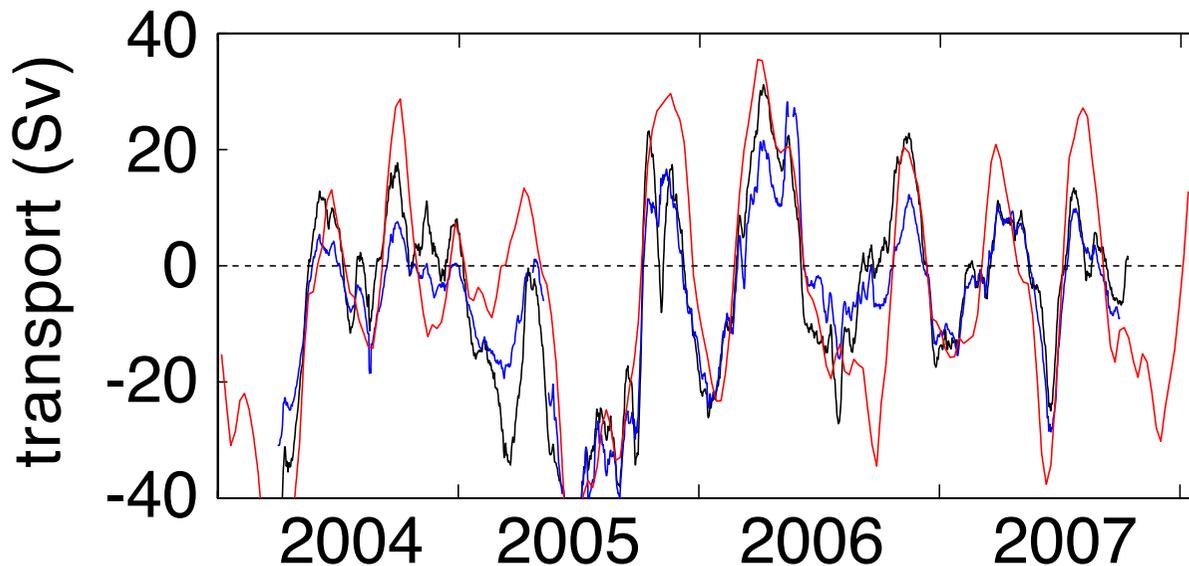
(the signal at 200 db is considered to avoid surface seasonal heating and cooling)

Is the transport between WB5 and MarWest recovered?

Geostrophic transport relative to bottom

- directly measured (T_ϕ)
- reconstructed from BC1 mode (T_{BC1})
- reconstructed from SSH and BC1 (T_{SSH})

std (Sv)	correlation		
	T_ϕ	T_{BC1}	T_{SSH}
17	1	0.91	0.80
24		1	0.83
20			1

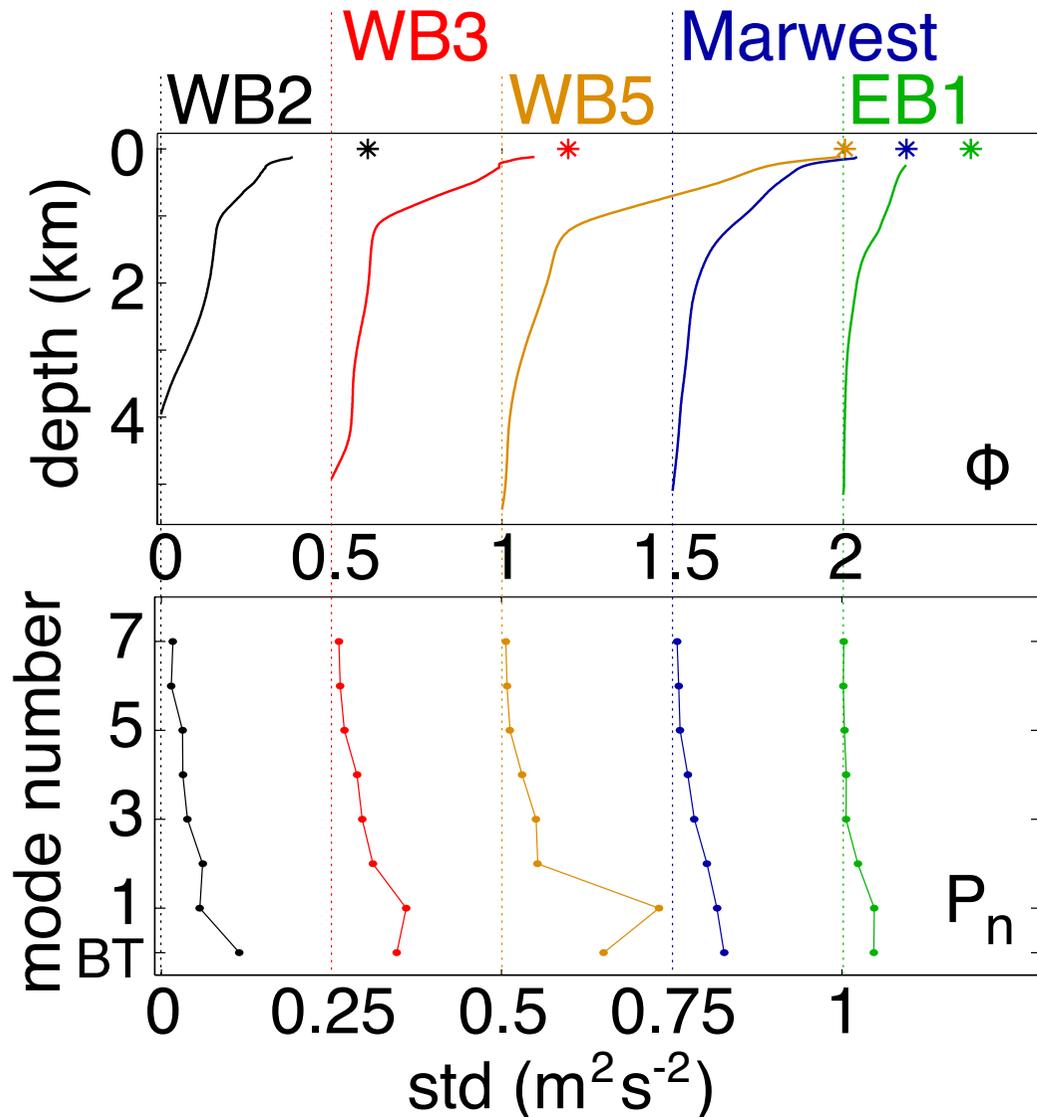


there's still good agreement:

of the variance in T_ϕ ,

- T_{BC1} recovers 83%
- T_{SSH} recovers 64%

How does mode decomposition vary across 26.5°N?



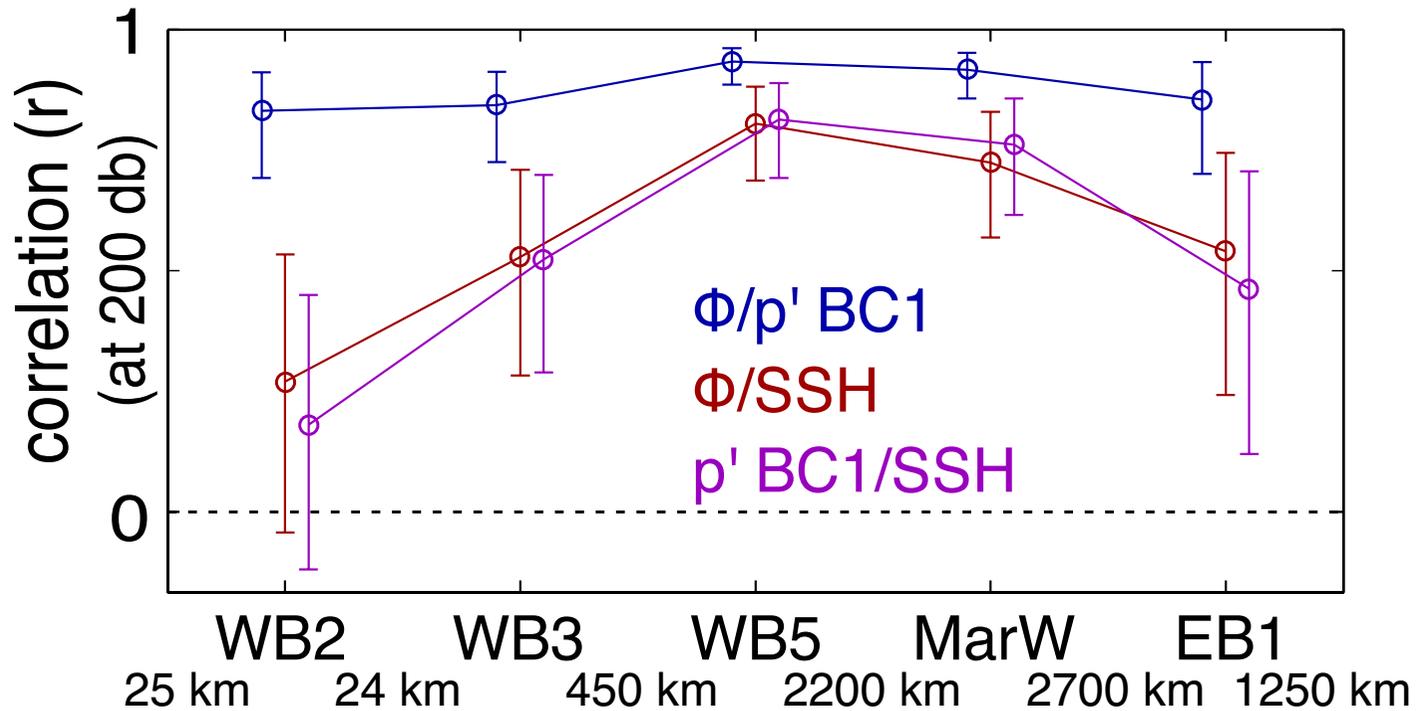
– std in upper 1000 is large and surface-intensified away from boundaries

– Surface intensification is weak at boundaries

– BC1 mode dominates away from boundaries

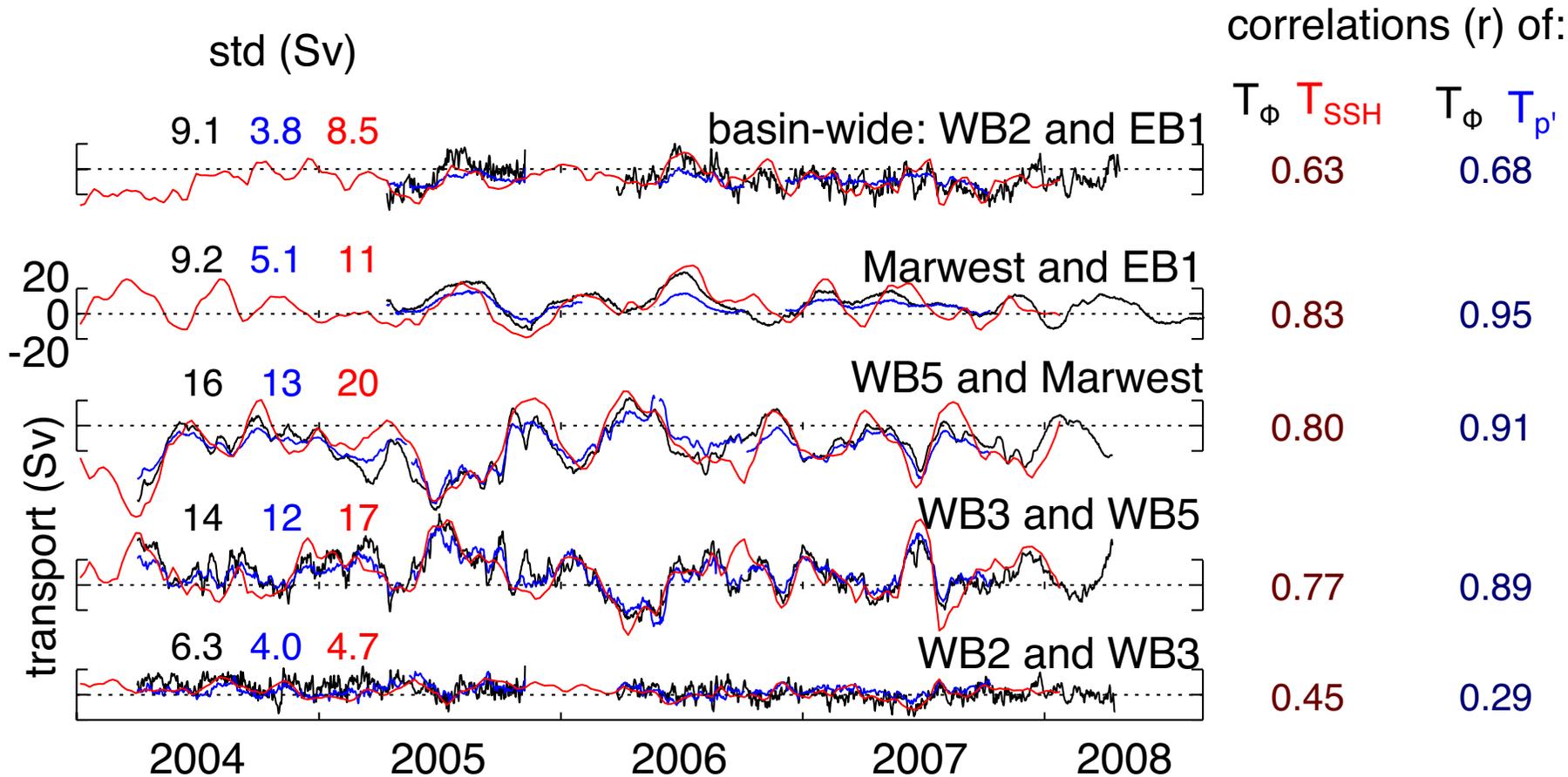
– WB2 and EB1 have different distribution of variance among modes

Correlations of the local signal across 26.5°N



- the **BC1-reconstruction** and Φ agree well at 200 db
- In the interior (WB5, Marwest), the **BC1-reconstruction** and **SSH** agree well with Φ
- At boundaries, although the **BC1-reconstruction** recovers the Φ signal that is correlated with **SSH**, neither is well-correlated with Φ .

Transports between moorings (T_ϕ , T_{BC1} , and T_{SSH})



- The local BC1 signal is also the dominant transport away from boundaries
- At boundaries, only 40% of the transport variance is recoverable from BC1 or SSH

Conclusions

We characterize the variance measured by density moorings in terms of the vertical mode structure.

- The first baroclinic mode dominates in ocean basin . . .
. . . but not within 25–50 km of W or 1000 km of E boundaries.
- Transport at the boundaries is poorly explained by the first baroclinic mode
. . . and so wave perturbations there do not have the simple form expected.
- This vertical structure offers a clear explanation of the limitations of recovering transport from SSH at this latitude.



Modal Decomposition Theory

A motionless reference state $\bar{\rho}(z)$ defines the perturbations, and its stratification N^2 defines $F(z)$ and $G(z)$

$$\frac{d}{dz} \left(\frac{1}{N^2} \frac{dF}{dz} \right) + \gamma^2 F = 0$$

and
$$G(z) = \frac{i\omega}{N^2(z)} \frac{dF(z)}{dz}$$

there are two vertical shapes (Gill, 1982; Wunsch and Stammer 1997):

$$u' = U(x, y, t) F(z)$$

$$v' = V(x, y, t) F(z)$$

$$p' = \rho_0 P(x, y, t) F(z)$$

$$w' = P(x, y, t) G(z)$$

$$\xi' = H(x, y, t) G(z)$$

The equivalent geopotential anomaly modes ϕ'_n are calculated with 1) hydrostatic balance to obtain ρ'_n , 2) vertical integration to obtain ϕ'_n , and 3) the Boussinesq approximation:

$$\Phi'_n = \int_0^{p_0} \delta'_n dp \approx \int_0^{p_0} \frac{-\rho'_n}{\bar{\rho}^2} dp = \int_0^{p_0} \frac{-1}{\bar{\rho}^2} \frac{\partial p'_n}{\partial z} dp = \dots = \frac{1}{\rho_0} p'_n$$