

Ice-Ocean Boundary Dynamics

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Outline



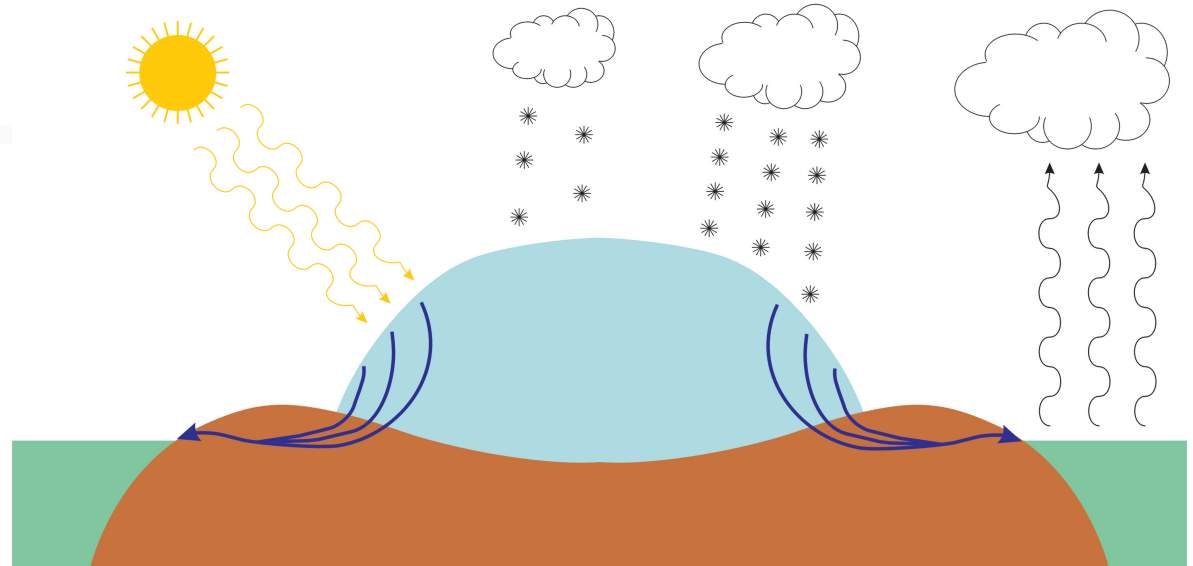
- I. Brief motivation
- II. Parameterisation of the turbulent boundary layer
- III. Buoyancy-driven boundary currents
- IV. Summary of knowns and unknowns



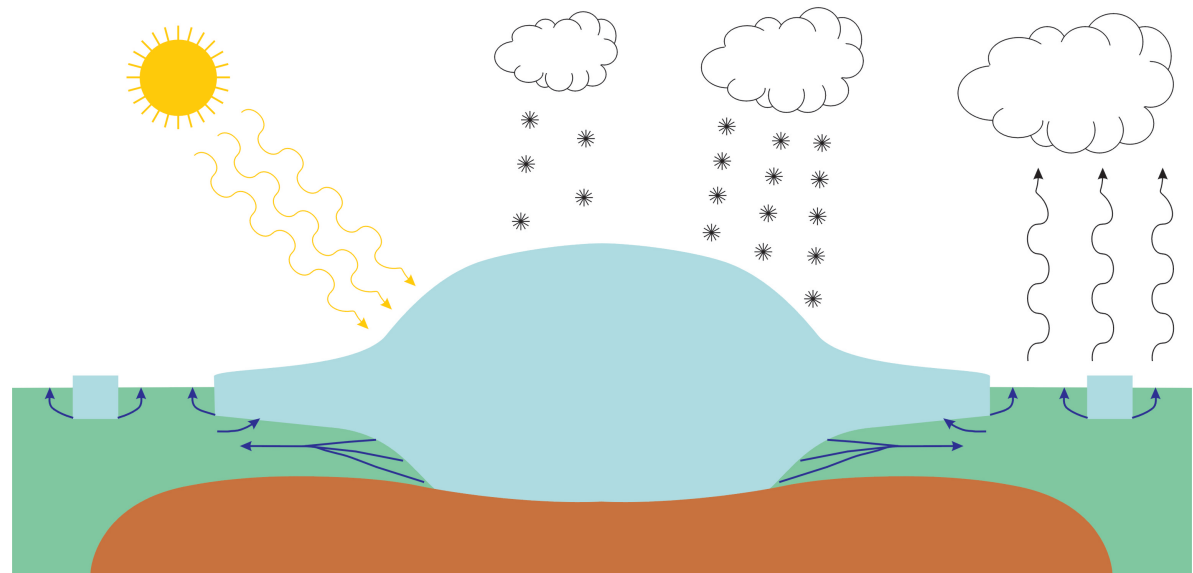
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Terrestrial ice sheet

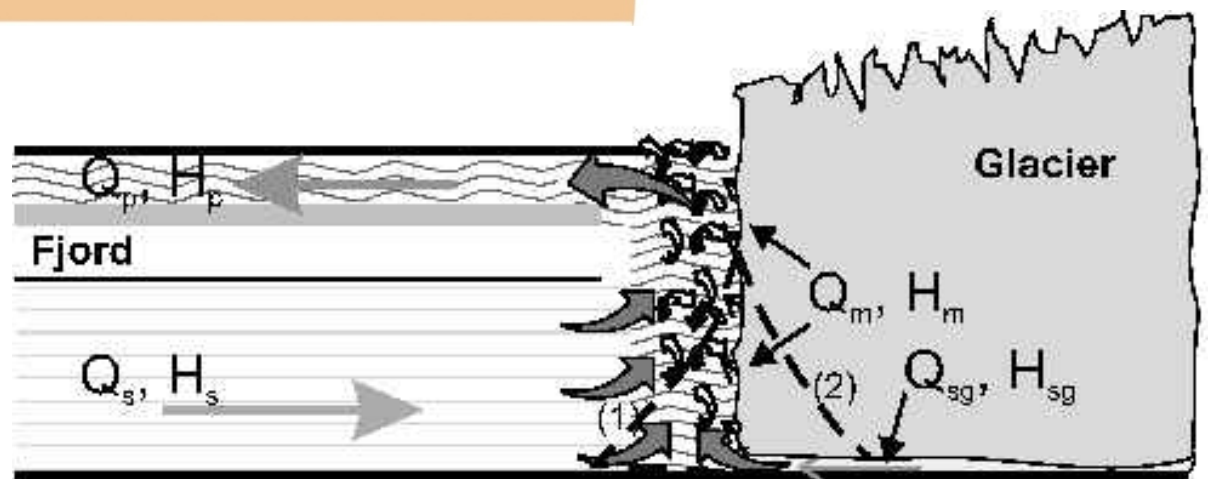
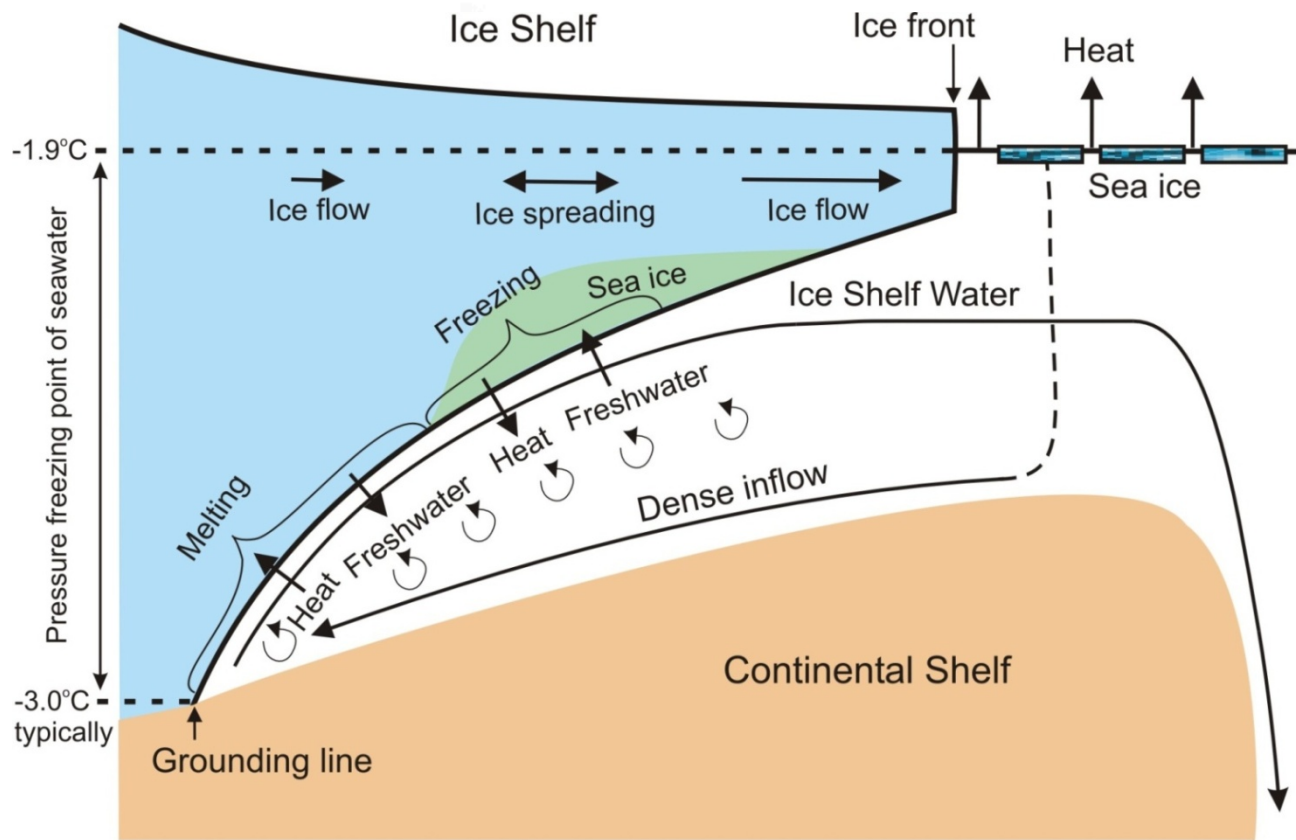


Marine ice sheet



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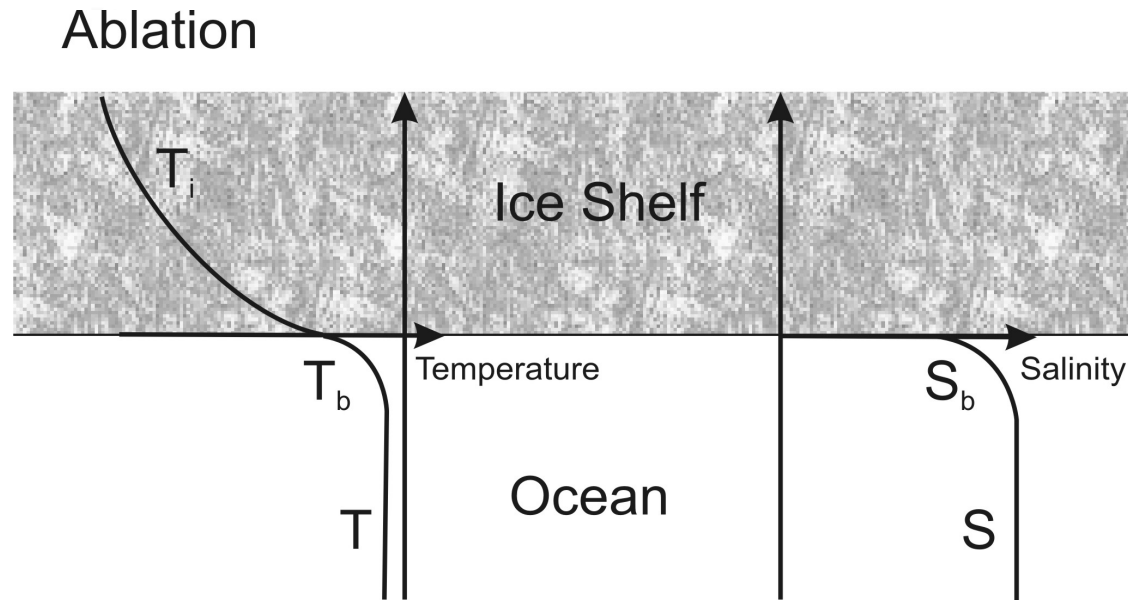
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Motyka et al. (2003)



Liquidus condition:

$$T_b = \lambda_1 S_b + \lambda_2 + \lambda_3 P_b$$

Heat balance:

$$\rho_i \dot{a}_b L = \rho_i c_i \kappa_i \left. \frac{\partial T_i}{\partial n} \right|_b - q_b^T$$

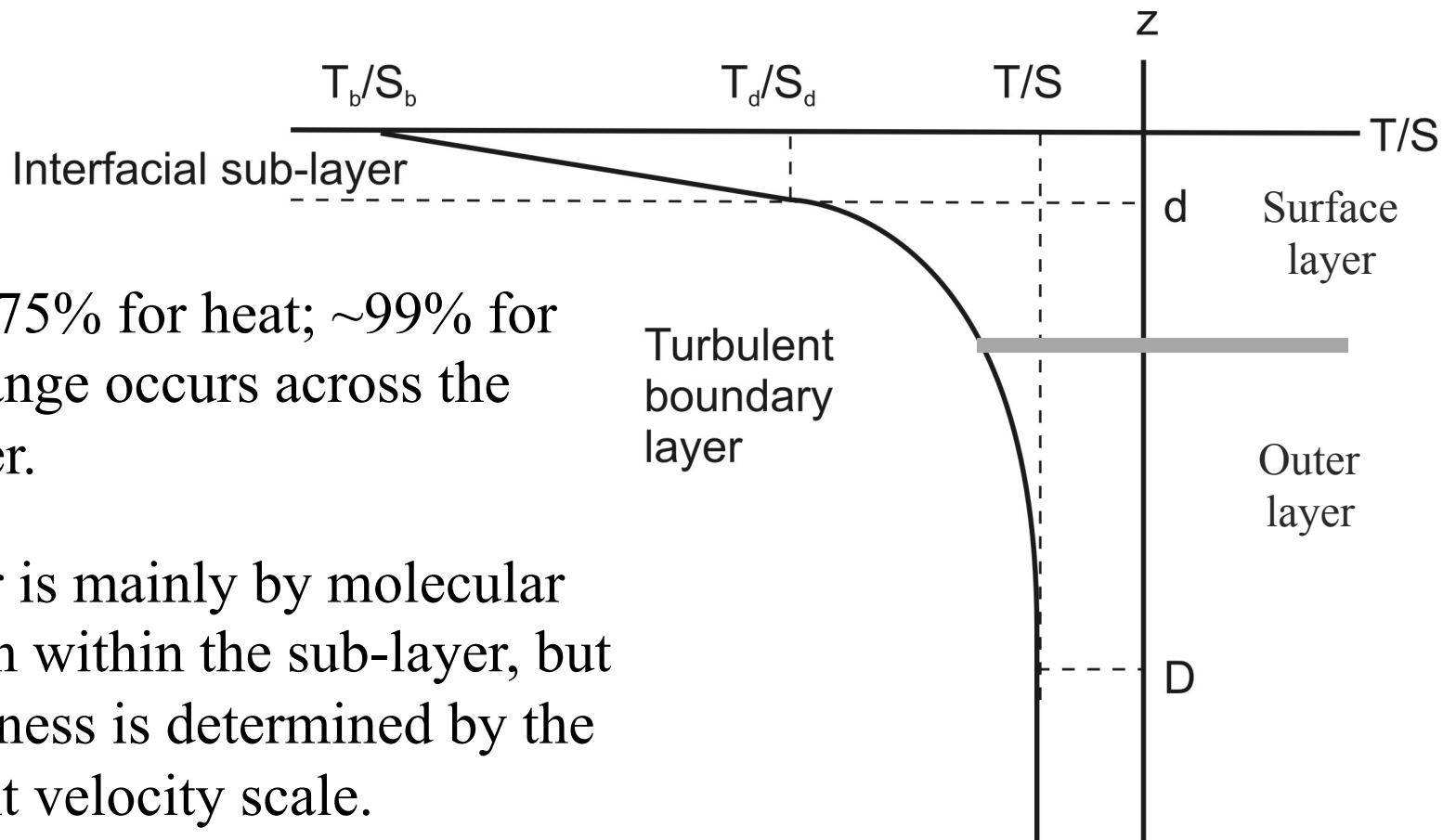
Salt balance:

$$\rho_i \dot{a}_b (S_b - S_i) = -q_b^S$$



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Most (~75% for heat; ~99% for salt) change occurs across the sub-layer.

Transfer is mainly by molecular diffusion within the sub-layer, but its thickness is determined by the turbulent velocity scale.

Interfacial shear stress:

$$\tau_b = \rho_w C_d U^2$$

Turbulent fluxes:

$$q_b^T = \rho_w c_w C_d^{1/2} U \Gamma_T (T_b - T)$$

$$q_b^S = \rho_w C_d^{1/2} U \Gamma_S (S_b - S)$$



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$$C_d$$



Almost entirely unknown for ice shelves/tidewater glaciers.

Typically assumed to be $\sim 10^{-3}$.

Similar to values estimated for sea ice and the seabed.

$$\Gamma_T, \Gamma_S$$

Almost entirely unknown for ice shelves/tidewater glaciers.

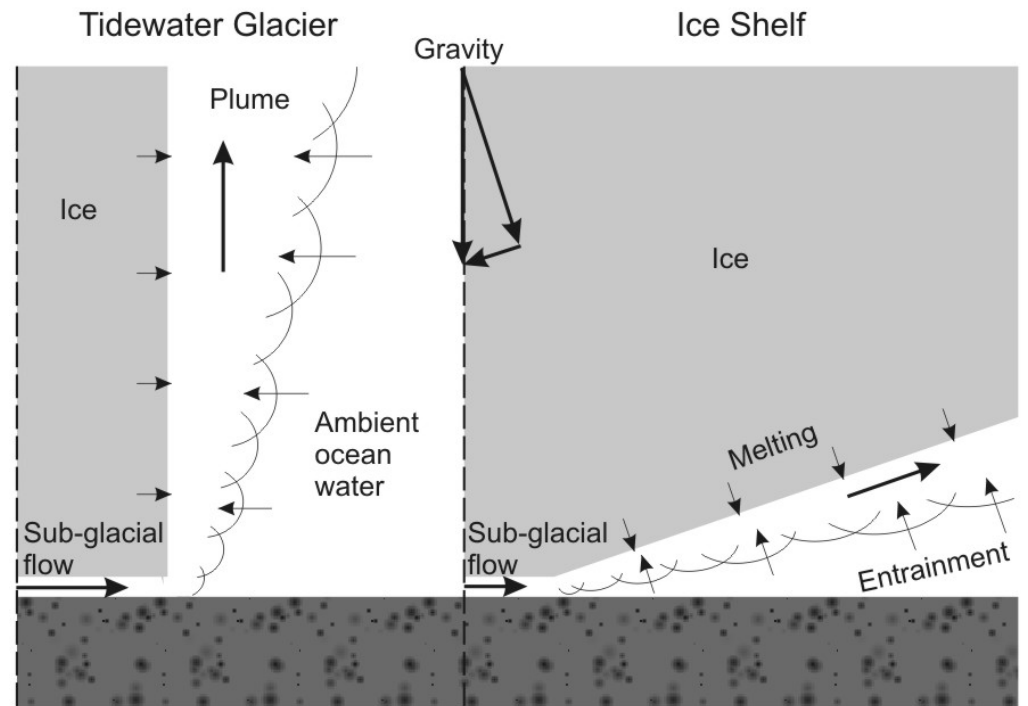
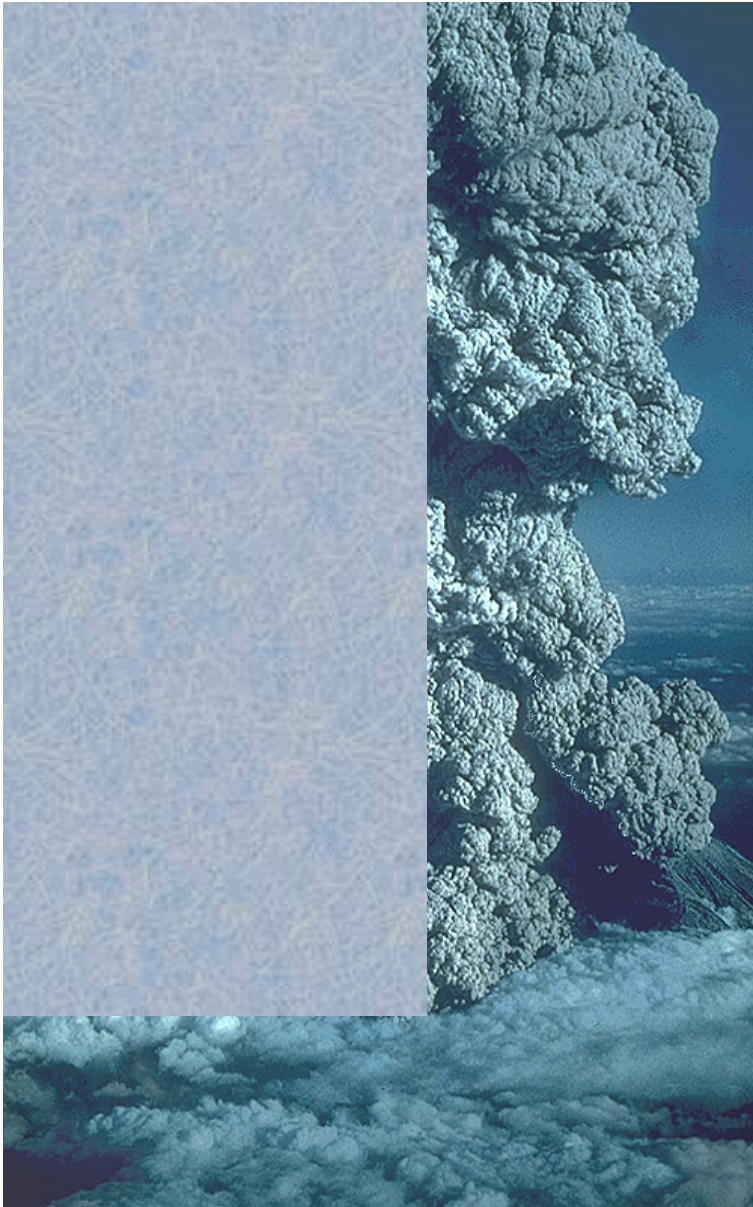
Expressions taken from laboratory experiments or sea ice observations.

Dominant role of the molecular processes in the interfacial sublayer suggests that values should be broadly applicable.



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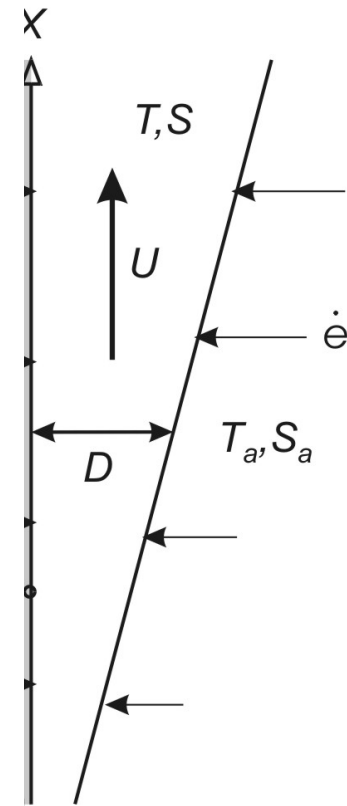
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$$\frac{d}{dX}(DU^2) = D\Delta\rho g \sin \alpha - C_d U^2$$

$$\frac{d}{dX}(DU) = \dot{e}$$

$$\frac{d}{dX}(DUT) = \dot{e}T_a$$

$$\frac{d}{dX}(DUS) = \dot{e}S_a$$



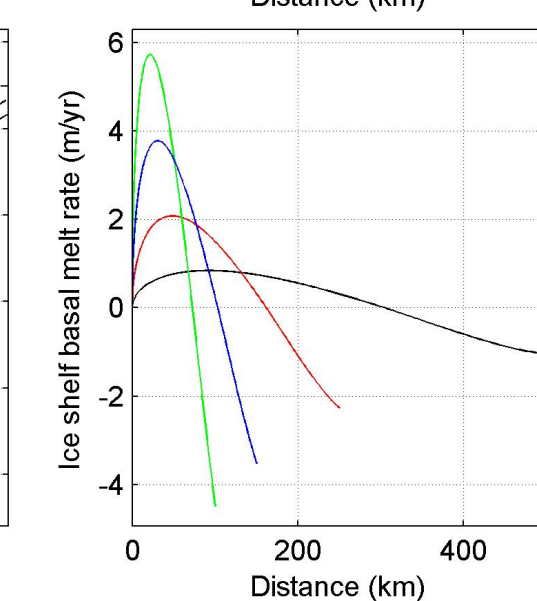
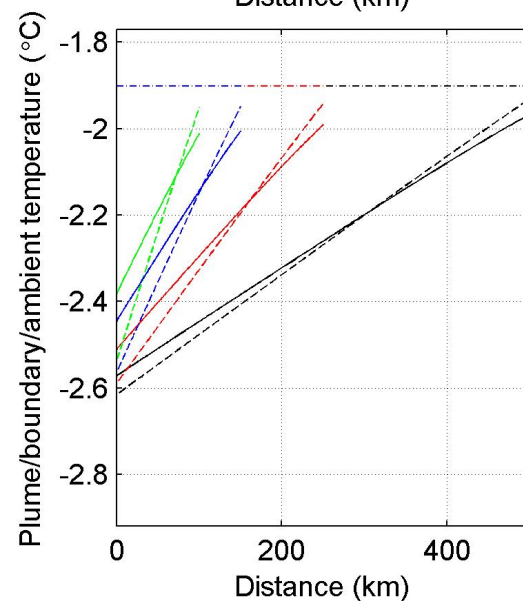
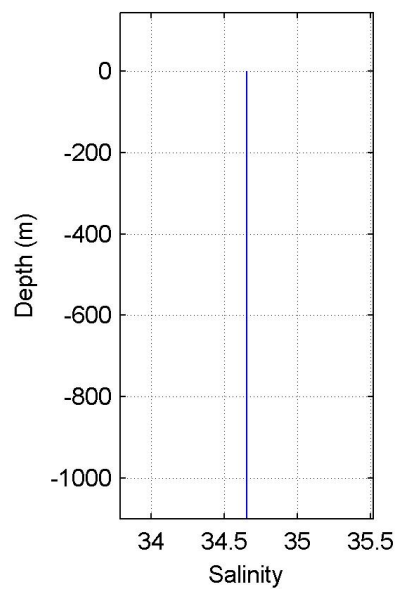
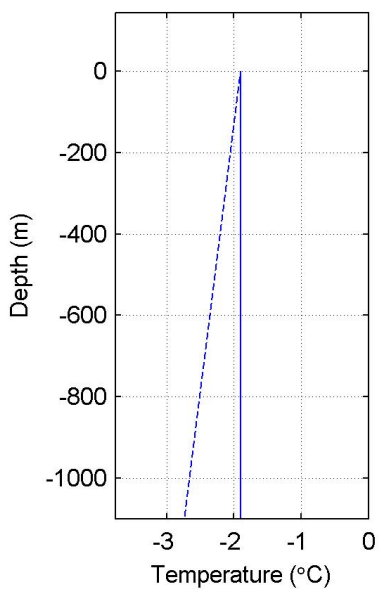
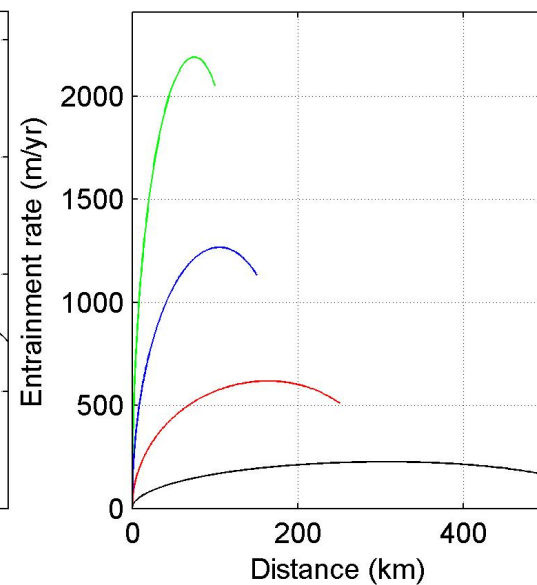
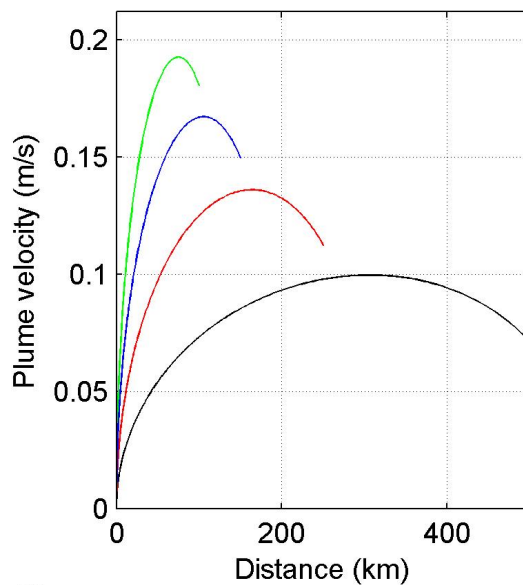
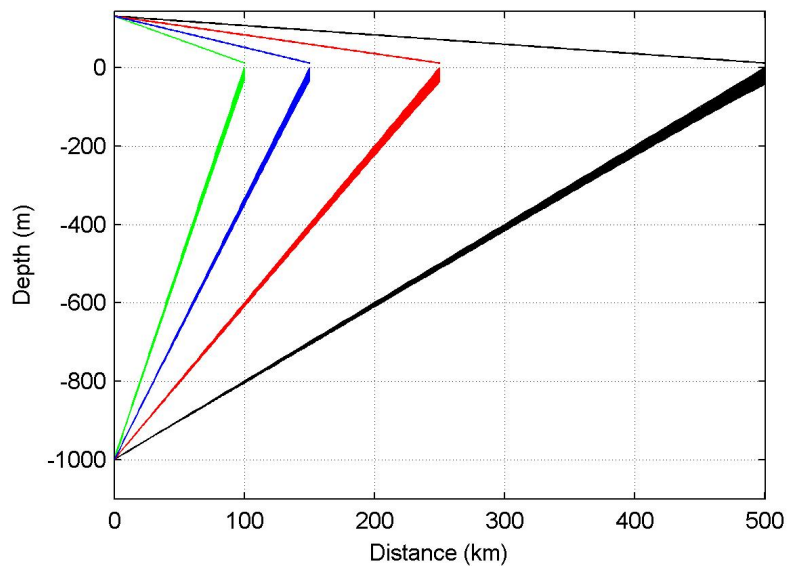
$$\Delta\rho = \beta_S(S_a - S) - \beta_T(T_a - T)$$

$$\dot{e} = E_0 U \sin \alpha$$



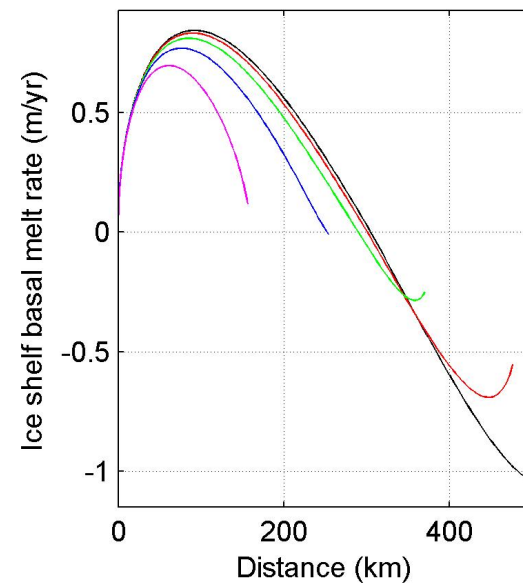
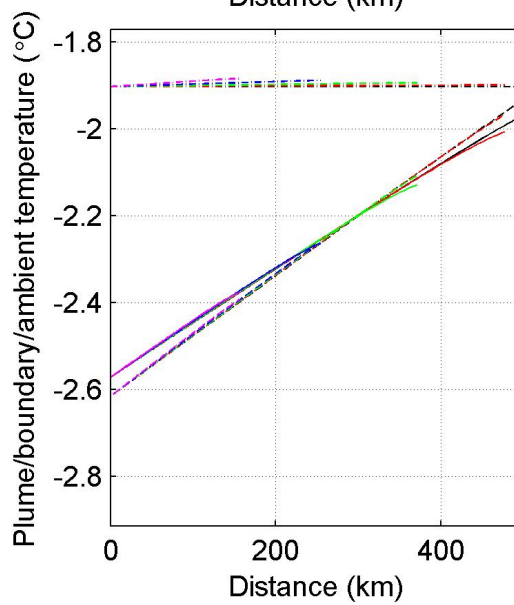
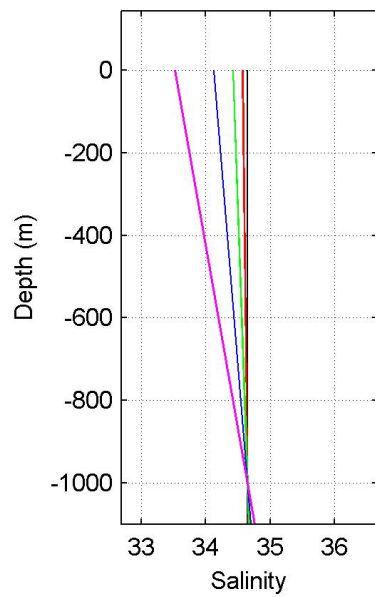
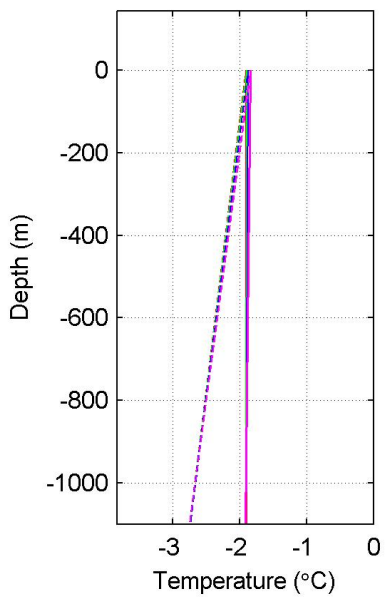
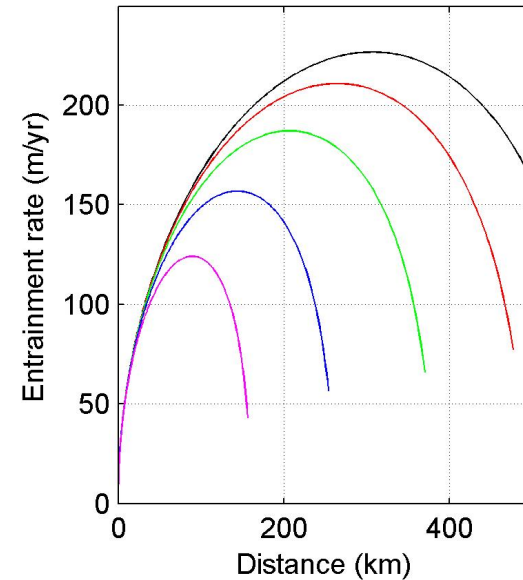
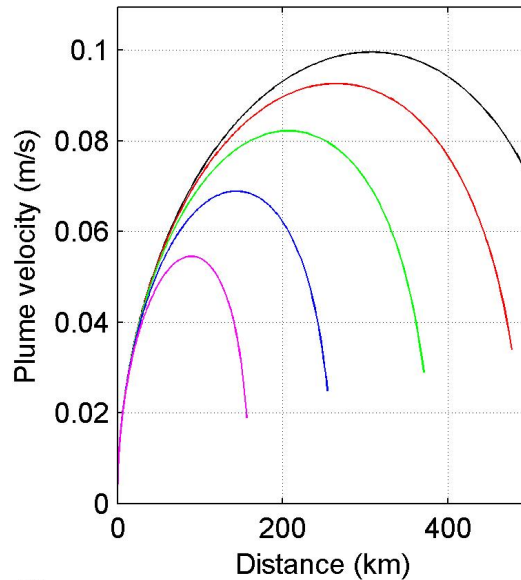
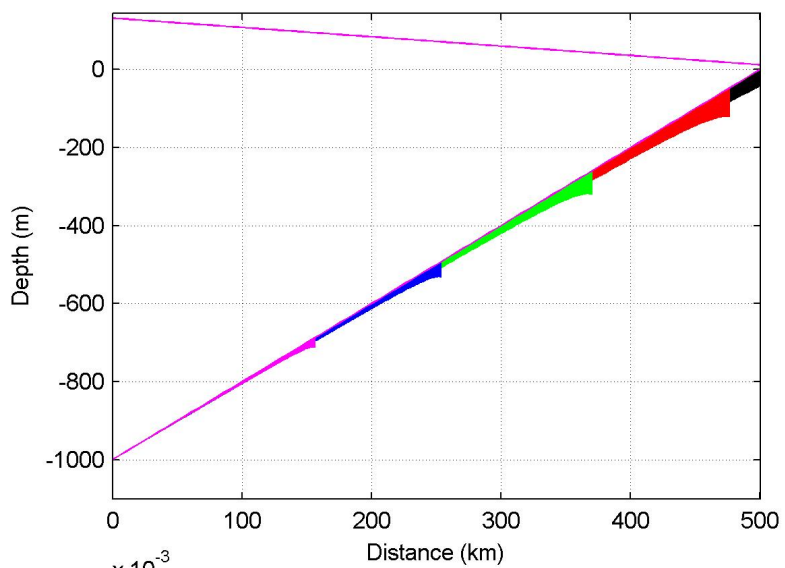
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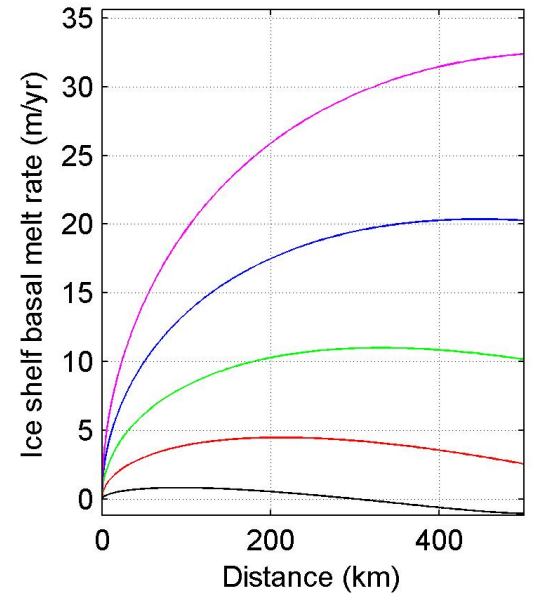
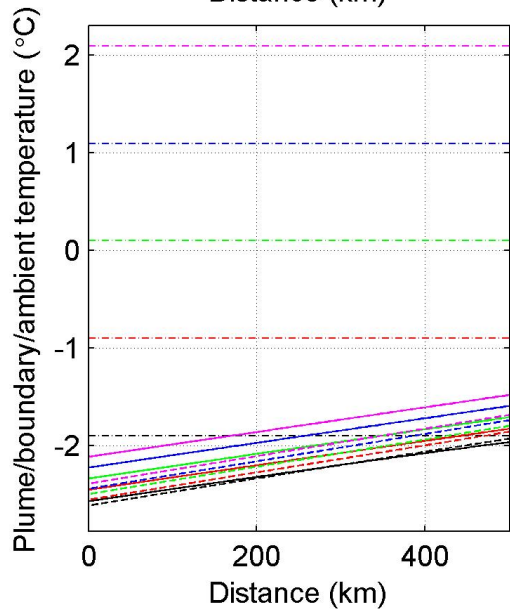
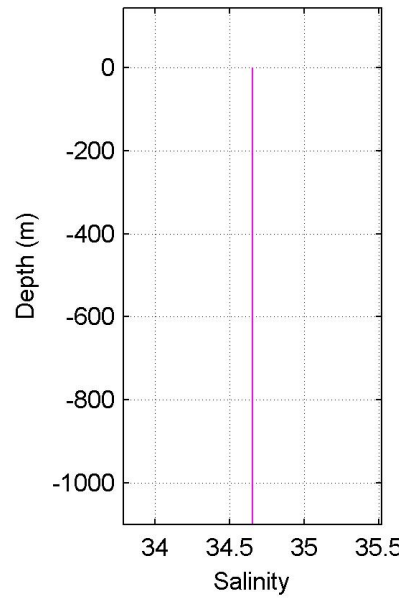
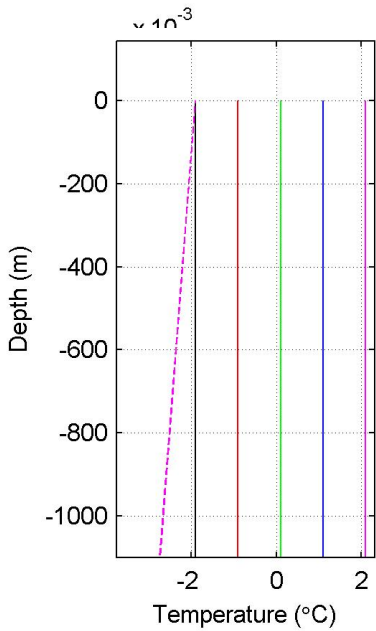
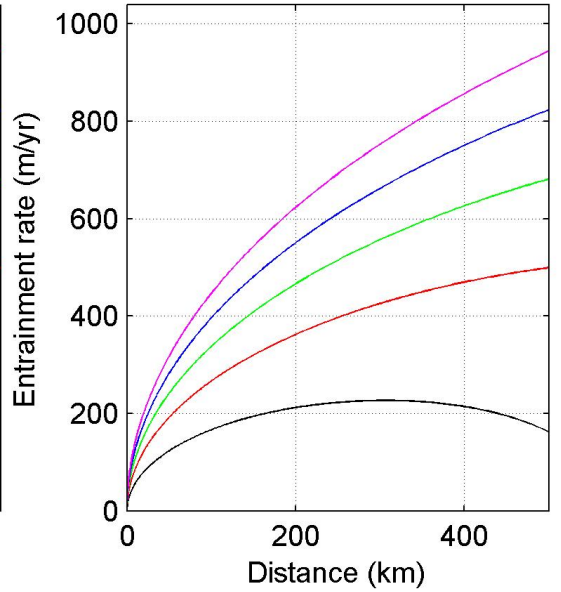
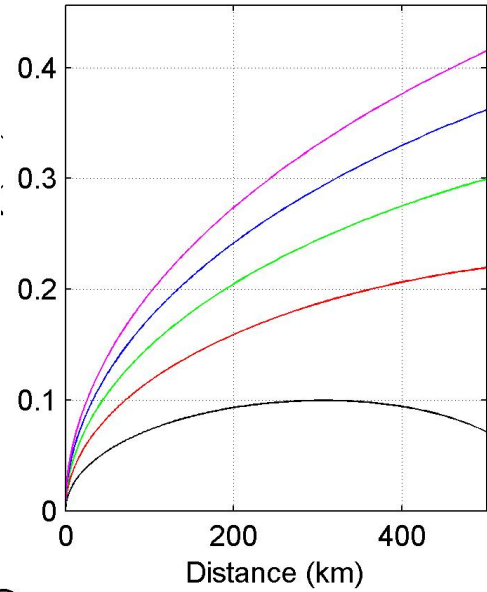
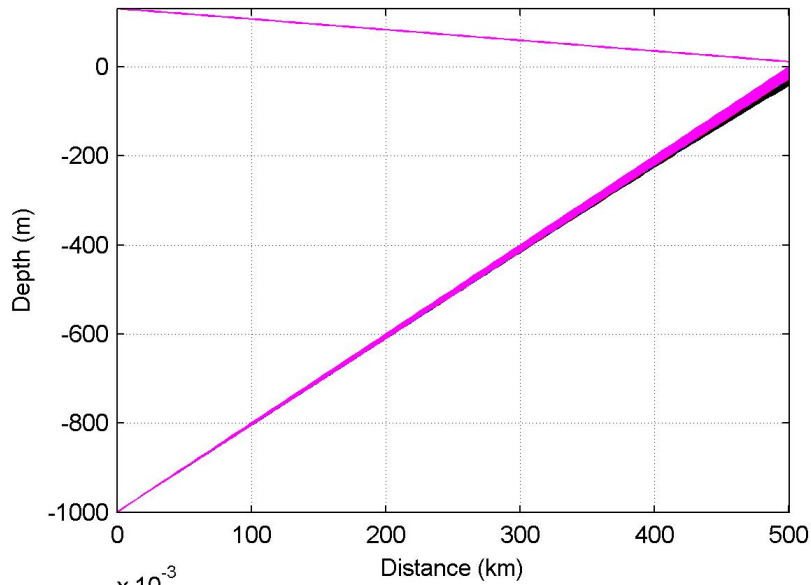
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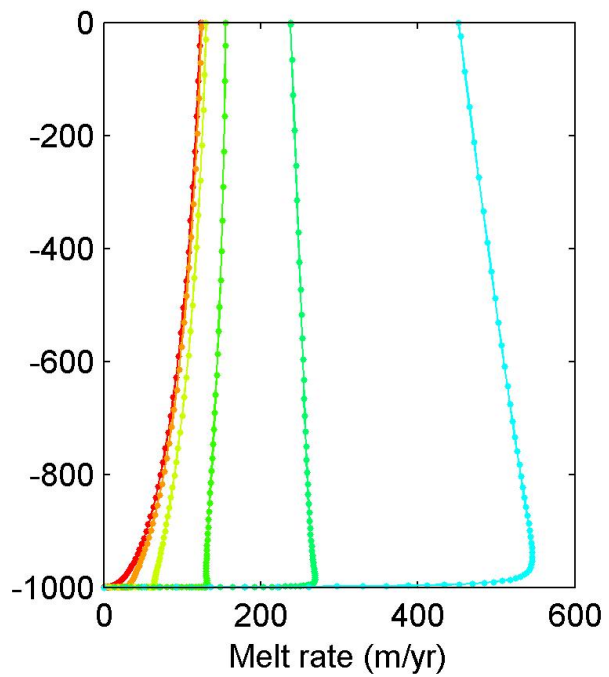
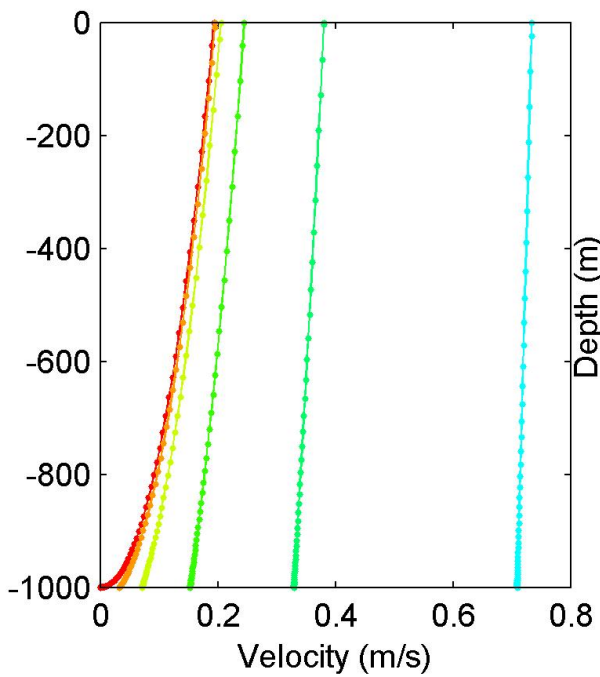
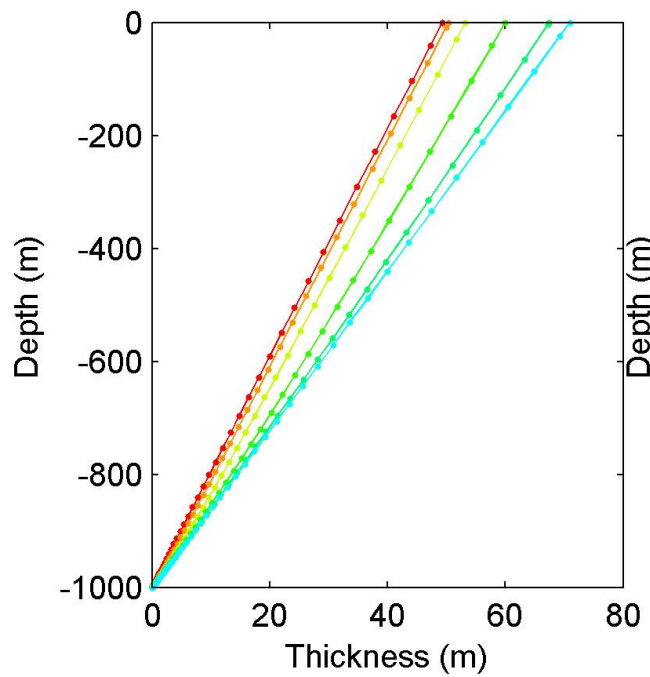
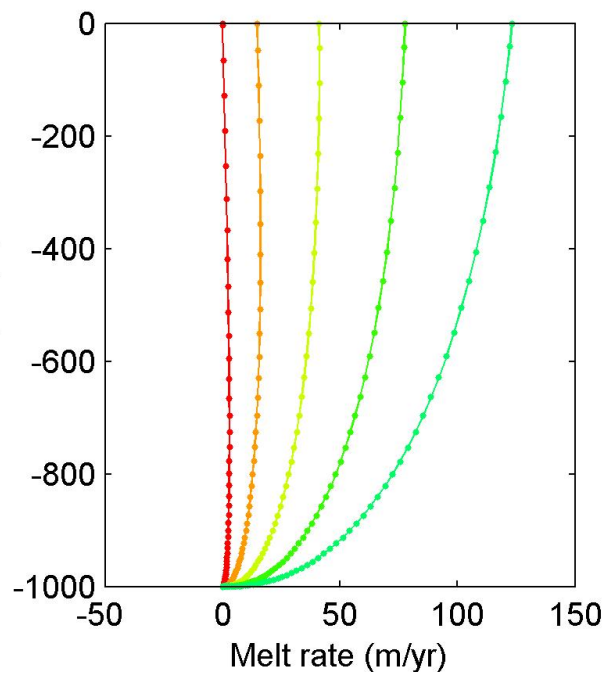
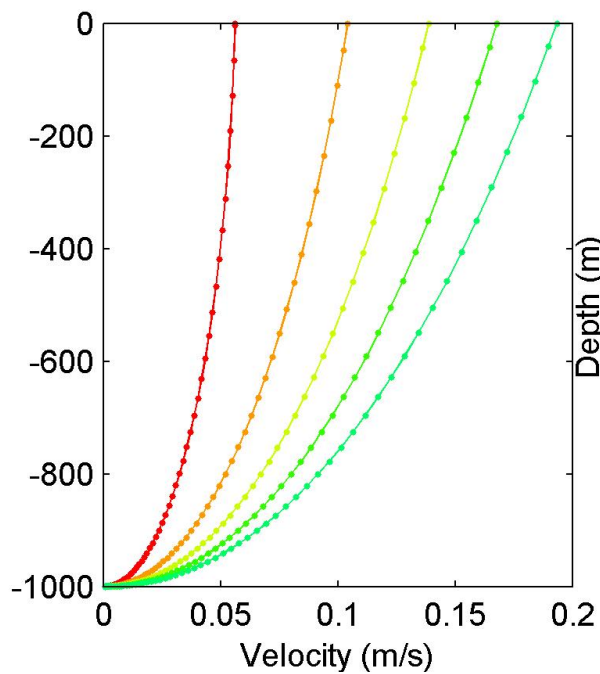
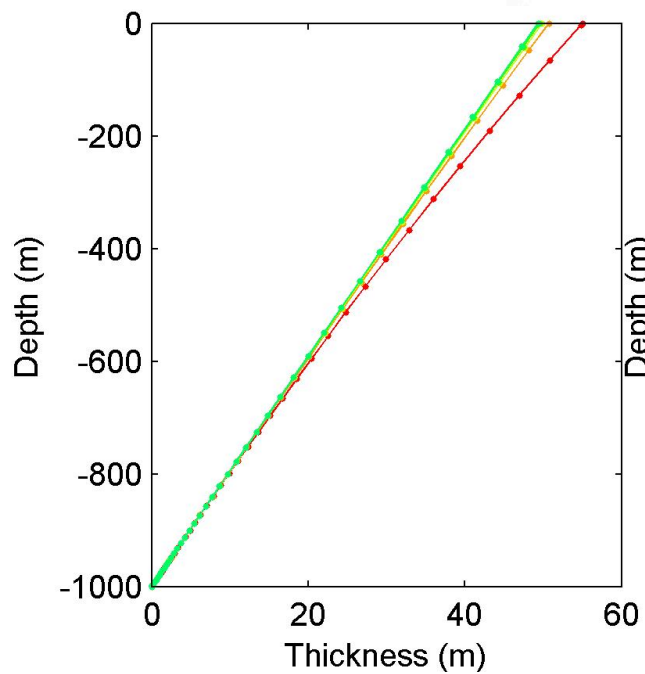
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If we assume melting has no impact on plume buoyancy, then in an unstratified environment there is a simple solution (if we also neglect the pressure dependence of the freezing point):

constant velocity,
$$U_0 = \left(\frac{B_0 \sin \alpha}{E_0 \sin \alpha + C_d} \right)^{1/3}$$

linearly increasing thickness,
$$D = E_0 \sin \alpha X$$

constant buoyancy flux,
$$B_0 = gD_0 U_0 \Delta \rho_0$$

constant thermal driving,
$$(T - T_f)_0 = \left(\frac{E_0 \sin \alpha}{E_0 \sin \alpha + C_d^{1/2} \Gamma_{\{TS\}}} \right) (T_a - T_{af})$$

and constant melt rate,
$$\dot{m}_0 = \frac{c C_d^{1/2} \Gamma_{\{TS\}}}{L'} U_0 (T - T_f)_0$$



The melt rate can therefore be written as:

$$\dot{m}_0 = \left(\frac{cC_d^{1/2}\Gamma_{\{TS\}}}{L'} \right) \left(\frac{\sin \alpha}{E_0 \sin \alpha + C_d} \right)^{1/3} \left(\frac{E_0 \sin \alpha}{E_0 \sin \alpha + C_d^{1/2}\Gamma_{\{TS\}}} \right) B_0^{1/3} (T_a - T_{af})$$

Physical constants

Geometrical factors

Cube root of buoyancy flux at grounding line

Ambient thermal driving

For a vertical ice face: $\sin \alpha = 1$

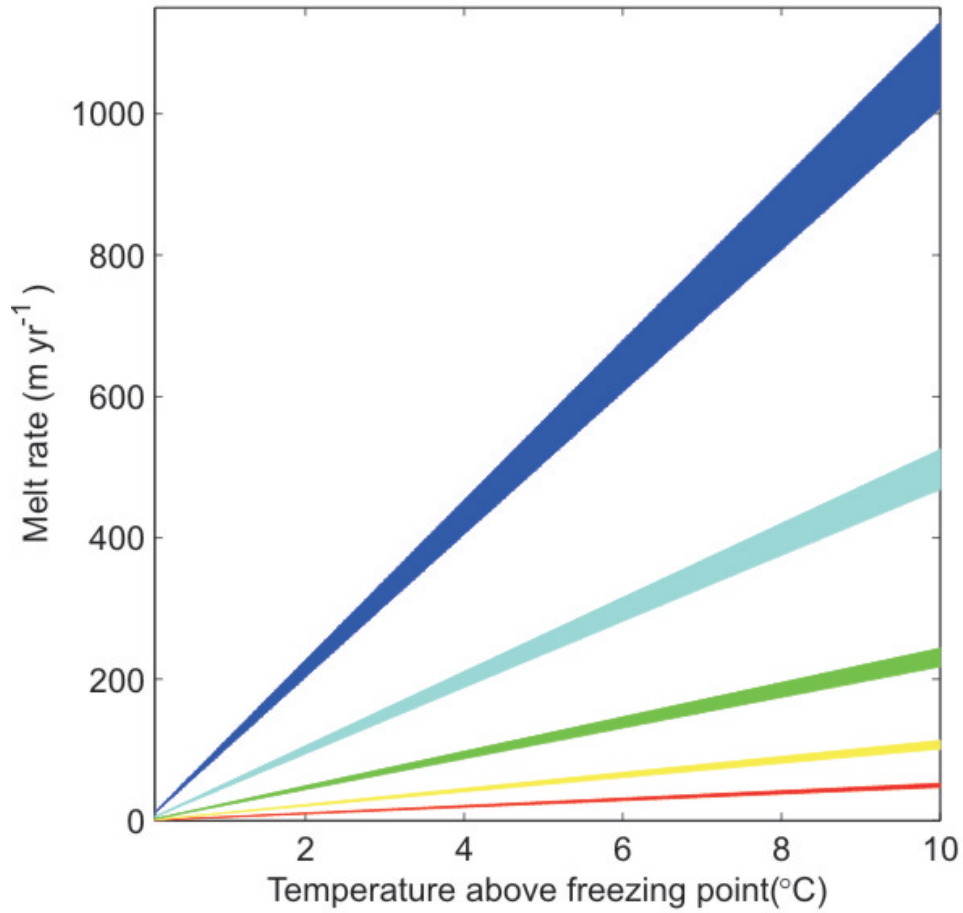
$$E_0 \gg C_d \gg C_d^{1/2}\Gamma_{\{TS\}}$$

$$\Rightarrow \text{Geometrical factors} \approx E_0^{-1/3}$$



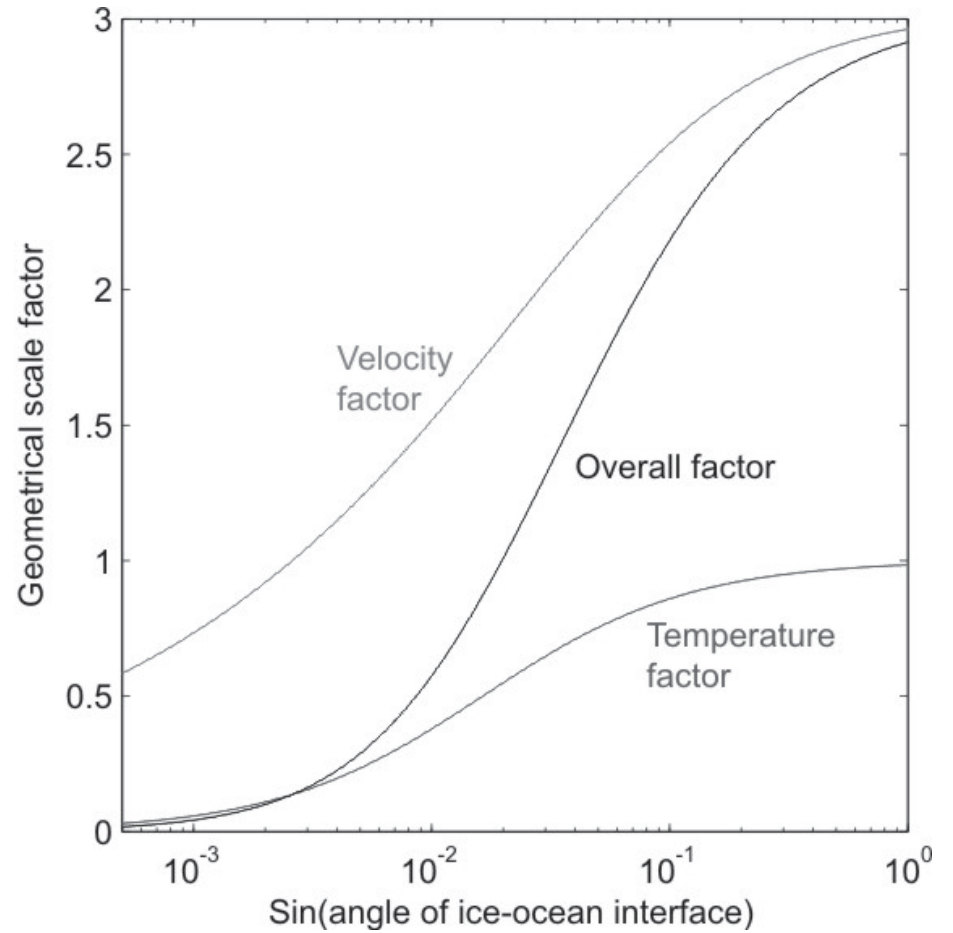
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Initial freshwater flux (m² s⁻¹):

— 5e-5 — 5e-4 — 5e-3 — 5e-2 — 5e-1



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Summary I

- Physics of the turbulent ice-ocean boundary layer beneath sea is fairly well understood.
- Parameterisations of turbulent transfer are well tested against observation, but those observations inevitably sample a limited range of conditions.
- Parameterisations should be readily transferable to ice shelves and tidewater glaciers, particularly the latter.
- There are very few observations to confirm or refute that, although data are now being collected beneath ice shelves.
- Sampling near a vertical ice face is challenging!



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Summary II

- Plume theory has been applied successfully to understand the fundamental behaviour of buoyancy-driven boundary currents.
- Once again, the theory should be more applicable to tidewater glaciers than ice shelves.
- Again there are no observations to support that claim.
- Recent work (Xu et al, 2012; Sciascia et al, in press, several posters presented here) have applied 2 and 3 dimensional ocean models to the vertical ice face problem.
- Very high resolution, non-hydrostatic codes are required to capture the processes.
- That makes the problem challenging for large-scale models!



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