



Context and Motivations

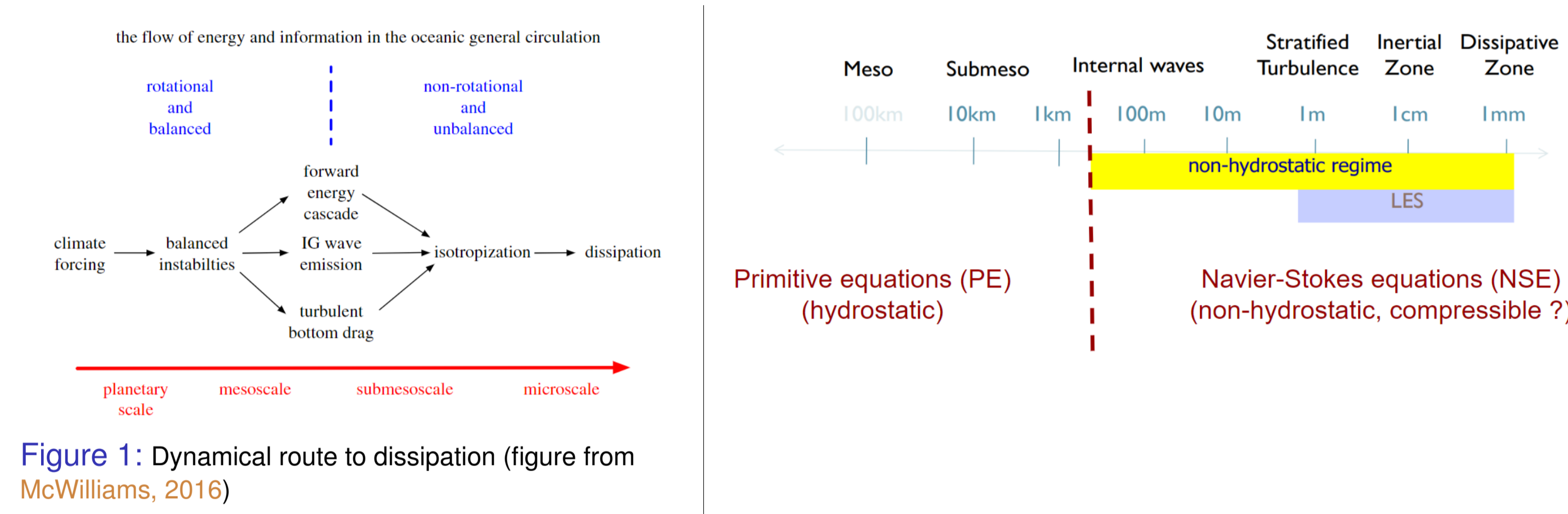


Figure 1: Dynamical route to dissipation (figure from McWilliams, 2016)

Focus of the project (funded from Feb. 2024 to Jan. 2028)

- ▷ Transition from hydrostatic to NH dynamics (occurs around a horizontal grid size of tens of meters)
- ▷ Representing both **downscaling** and **upscaling** across the submesoscale currents (SMCs) (feedback of NH processes on water-mass mixing and ocean circulation)

Range of scales of interest: submesoscale currents (SMCs)

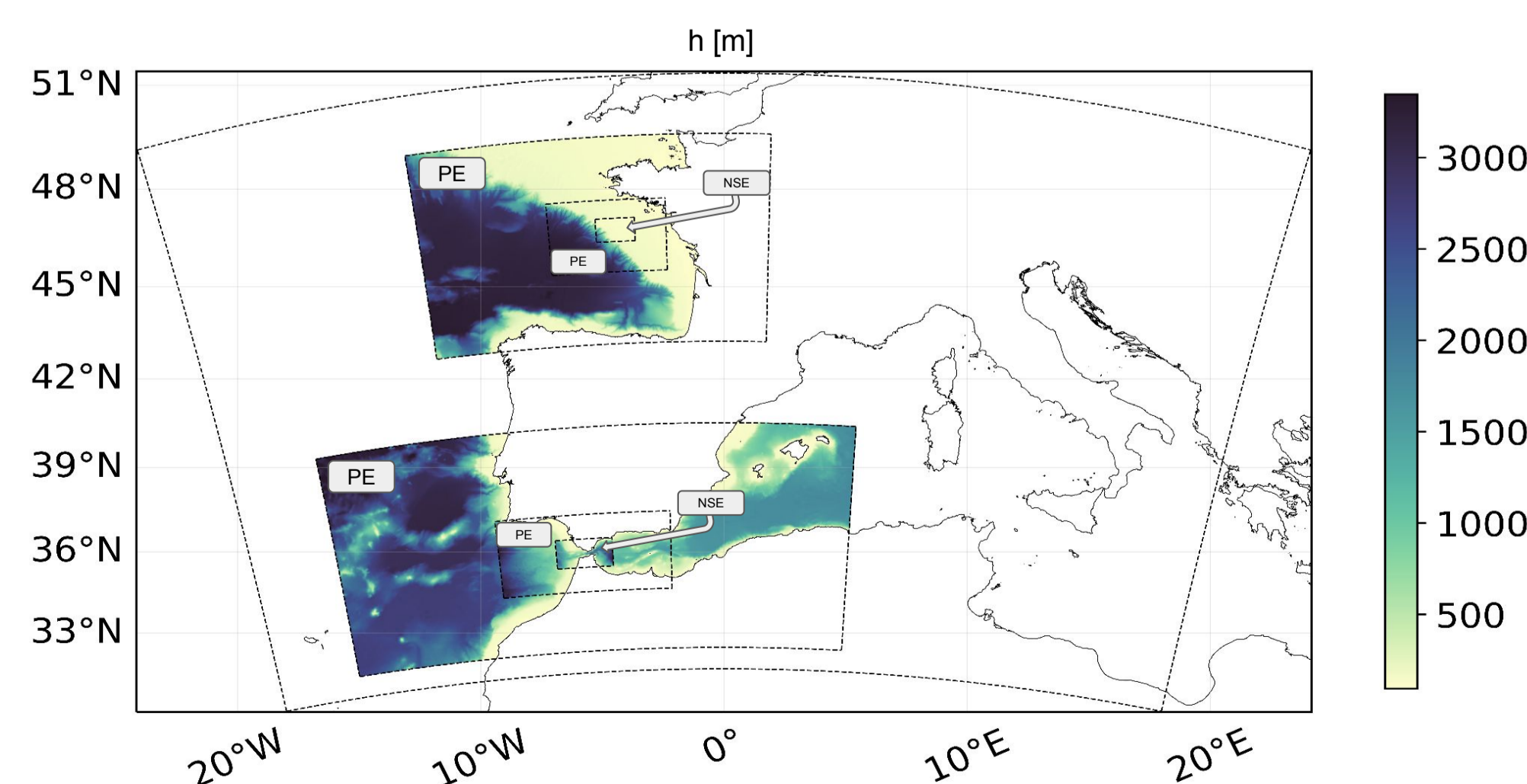
O(0.1 – 10 km/0.01 - 1 km) in the horizontal/vertical; O(days/hours) in time

“Large scales”: largest scales of the SMC range \equiv hydrostatically balanced dynamics
→ represented by most regional ocean forecasting systems

“Fine scales”: finest scales of the SMC range \equiv non-hydrostatic (NH) dynamics
→ not accessible for PE models

Objectives

- Provide robust and efficient numerical algorithms allowing a **multiscale modeling strategy based on block-structured mesh refinement with local adaptation of model equations, numerics and physics** in selected areas of interest.
- Assessment of developments through simulation of **fine-scale NH processes and their feedback to larger scales** within the Mediterranean / NE Atlantic dynamical continuum
- Comparing simulations with **in-situ observations** at fine scales available from Ifremer & SHOM in the Bay of Biscay and Gibraltar strait (towed Moving Vessel Profiler (MVP) CTD casts, moorings, drifters, Argo floats, eOdyn, etc)



⇒ Acquire new numerical perspectives for fine-scale dynamics and its feedback on the large scale.

Existing building blocks in the CROCO modeling system

CROCO: free-surface, structured grid oceanic model jointly developed by CNRS, Ifremer, Inria, IRD, SHOM and UT3



→ **Primitive Equation solver** (expanded on the basis of ROMS-AGRIF/UCLA)

Finite diff./vol. approach: C-grid staggering in the horizontal and **Lorenz-grid** in the vertical
Time-stepping: Leap-Frog-Adams-Mouton predictor corrector with split-explicit treatment of nonlinear free-surface
Momentum advection (conservative form) Linear schemes of order 2 to 6, WENO5z, TVD

Tracer advection
Horizontal: Linear schemes of order 2 to 6, WENO5z
Vertical: WENO5z, Compact4 (a.k.a. splines reconstruction)
Vertical coordinate: Quasi-Eulerian (terrain-following)

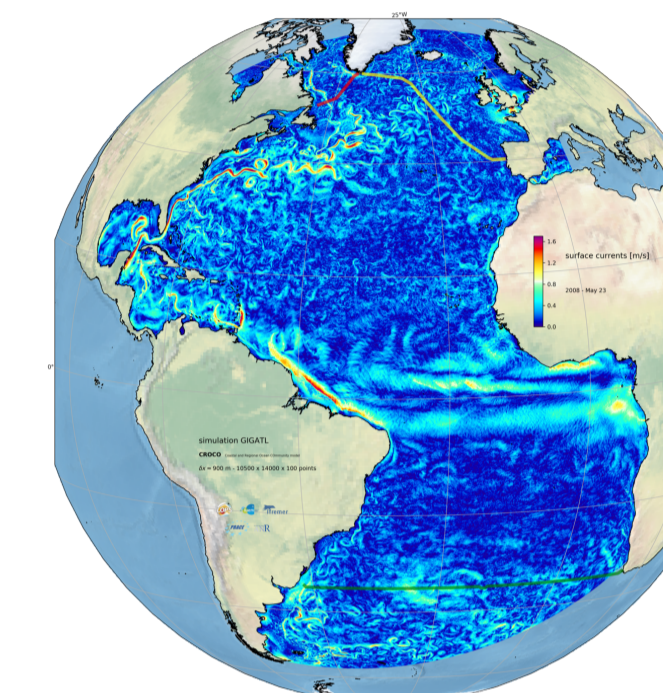


Figure 2: GIGATL Atlantic Ocean numerical simulations performed with CROCO (Gula et al., 2021)

→ **Navier-Stokes Equation solver**

Governing equations: compressible non-hydrostatic free-surface equations
→ retain fast acoustic waves and associated numerical stiffness

⇒ Design of a 2-mode, split-explicit time scheme for compressible NH ocean models (Auclair et al., 2024).

⇒ The fast *kernel* is 3d and no longer just barotropic.
⇒ Implemented as an overlay of the standard PE solver

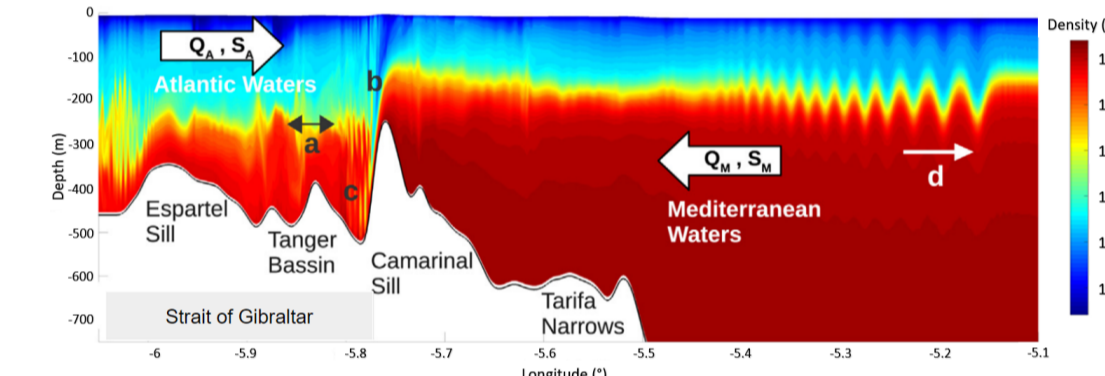
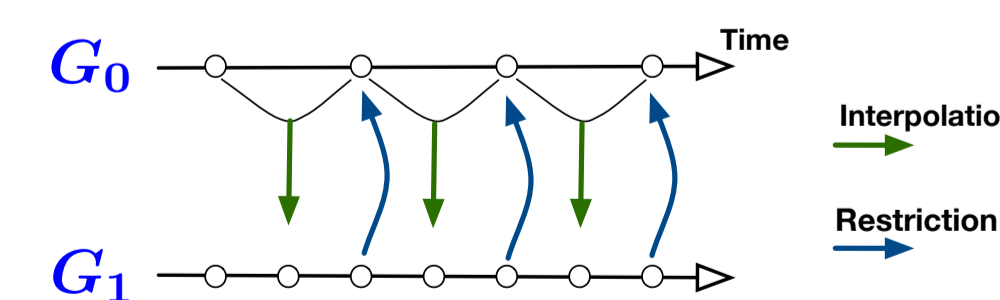


Figure 3: Illustration of fine-scale processes in the Strait of Gibraltar simulated using the CROCO NSE kernel in a 2DV configuration with a 50 m resolution. The visible processes include: (a) Small amplitude internal wave, (b) Hydraulic Jump, (c) KH instabilities, (d) Internal solitary waves. (Source: Hilt et al. (2020))

→ **2-way nesting capability**

- Full two-way coupling [Debreu et al., 2012]
→ solution unaffected by nesting when the refinement coefficient is one.
- Local space and time refinement (unlimited number of grids)
- Fully conservative (volume and tracer via refluxing)
- Numerical schemes and SGS parameterizations can differ from one level of resolution to another
- Implemented via the AGRIF library <http://agrif.imag.fr/>



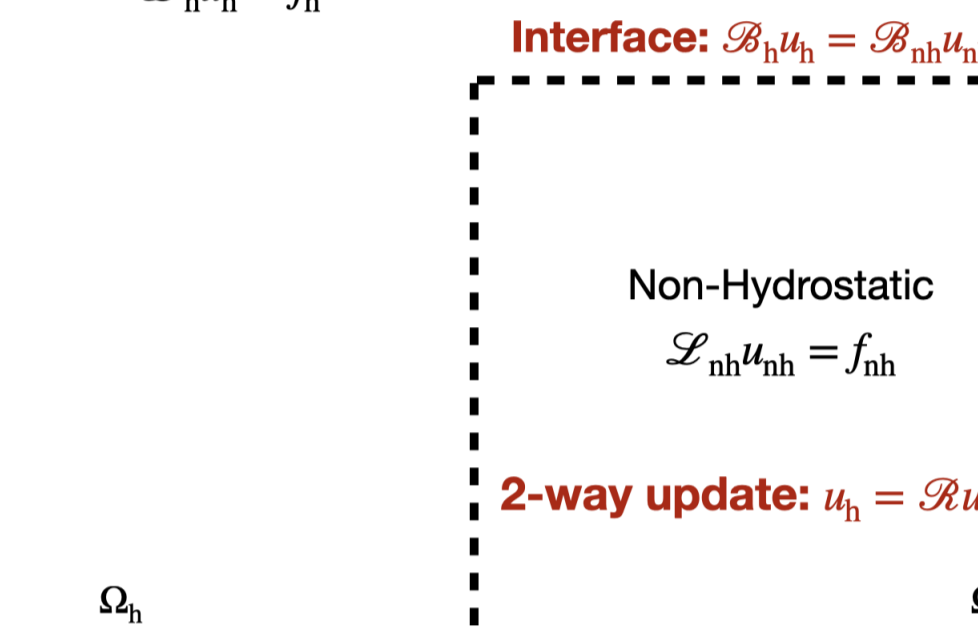
Features to be consolidated and/or developed for MOTIONS purposes

- Nesting of a high-resolution NSE zoom in a low-resolution PE grid with 2-way interactions
- Efficient numerical treatment of the fast acoustic/gravity waves for the NSE kernel (e.g. by considering only hydrostatic surface gravity waves)
- Deployment of large realistic simulations and associated HPC aspects (I/Os, pre/post processing, take advantage of the GPU version of PE/NSE solvers, ...)
- Subgrid-scale modeling for the O(10 m) resolution NSE zooms (Implicit vs explicit vs mixed explicit/implicit subgrid-scale-modeling)

Coupling hydrostatic and NH models

Problem: Find the interface operators \mathcal{B}_h and \mathcal{B}_{nh} (1-way) and the update operator \mathcal{R} (2-way)

Hydrostatic
 $\mathcal{L}_h u_h = f_h$



Non-Hydrostatic
 $\mathcal{L}_{nh} u_{nh} = f_{nh}$

2-way update: $u_{nh} = \mathcal{R} u_{nh}$

Objective: coupling methodology aiming at **matching vertical modes** combined with a **Perfectly Matched Layer (PML)** (Hu et al., 2008) (P. Lozano's thesis work)

Non-hydrostatic model
 $\rho_0 \partial_t u = -\partial_x p'$
 $\rho_0 \partial_t w = -\partial_z p' - g \rho'$
 $\partial_x u + \partial_z w = 0$
 $\partial_t \rho' = \frac{N^2}{g} \rho_0 w$

Boundary conditions

$w(x, z = -H, t) = 0$
 $\partial_x p'(x, z = 0, t) = \rho_0 g w(x, z = 0, t)$

→ IW dispersion relation:

$$\Omega^2 = \frac{N^2 k_x^2}{k_x^2 + k_z^2} < N^2$$

Associated Sturm-Liouville eigenvalue problem

$$W''(z) + \lambda (N^2(z) - \Omega^2) W(z) = 0$$

$$W(-H) = 0$$

$$W'(0) - g \lambda W(0) = 0$$

Expand the variables in the eigenfunctions to obtain a set of uncoupled systems

$$\partial_t u_q + \frac{g}{\lambda_q} \partial_x u_q = 0$$

$$\partial_t \rho_q + \frac{1}{g} (N^2 + \Omega^2) \partial_x u_q = 0$$

→ achieve coupling by matching the coefficients of each mode

Case study: coupling KdV/transport

• We want to solve:

$$\frac{\partial u}{\partial t} + c \frac{\partial u}{\partial x} + \epsilon \frac{\partial^3 u}{\partial x^3} = 0$$

$$\frac{\partial u}{\partial t} + c \frac{\partial u}{\partial x} = 0$$

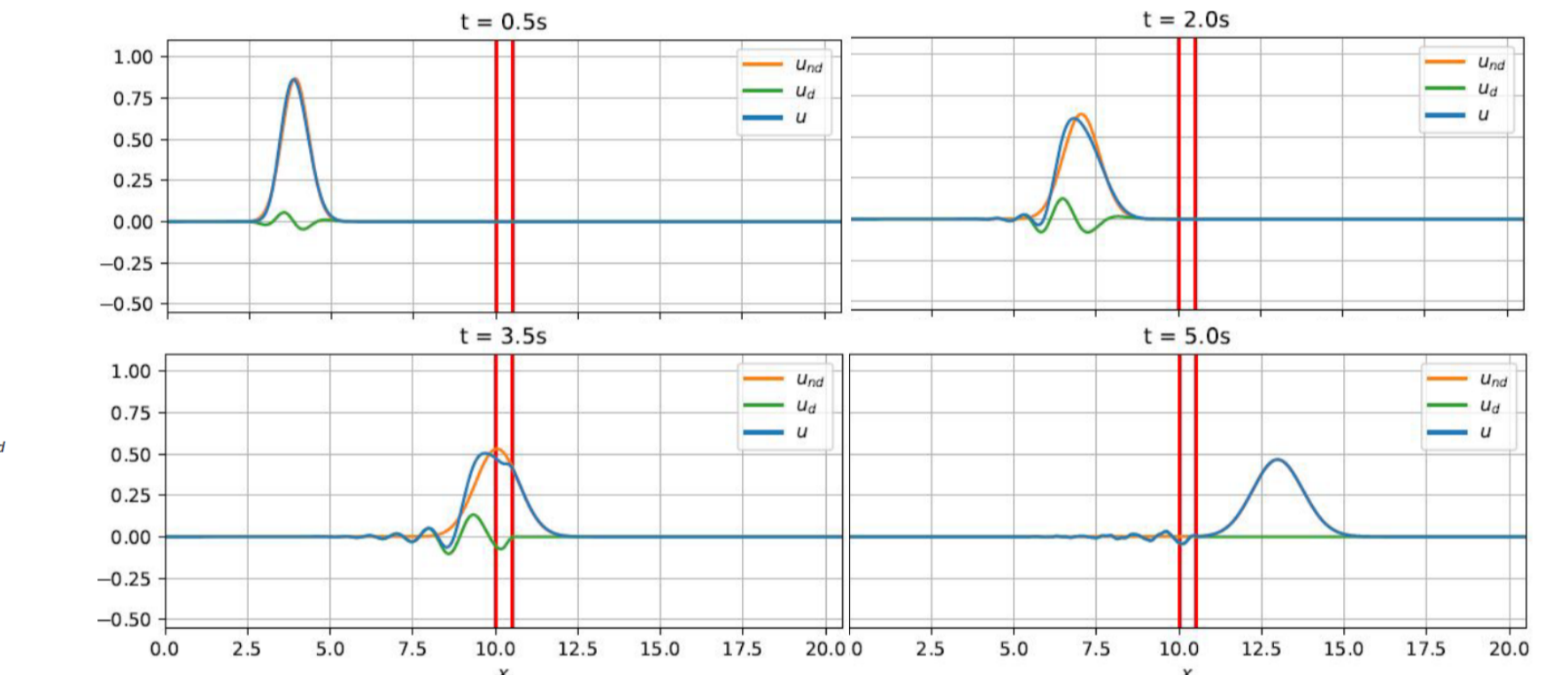
• Let $u = u_d + u_{nd}$ such that:

$$\begin{cases} \frac{\partial u_d}{\partial t} + c \frac{\partial u_d}{\partial x} = 0 & \text{in } \Omega_d \cup \Omega_{nd} \\ u_d = 0 & \text{in } \Omega_{nd} \end{cases} \Rightarrow \begin{cases} \frac{\partial u_d}{\partial t} + c \frac{\partial u_d}{\partial x} + \epsilon \frac{\partial^3 u_d}{\partial x^3} = -\epsilon \frac{\partial^3 u_{nd}}{\partial x^3} & \text{in } \Omega_d \\ \frac{\partial u_{nd}}{\partial t} + c \frac{\partial u_{nd}}{\partial x} = 0 & \text{in } \Omega_{nd} \end{cases}$$

• Add a PML to attenuate u_d exiting from Ω_d .

$$\frac{\partial u}{\partial t} + c \frac{\partial u}{\partial x} + \epsilon \frac{\partial^3 u}{\partial x^3} = 0$$

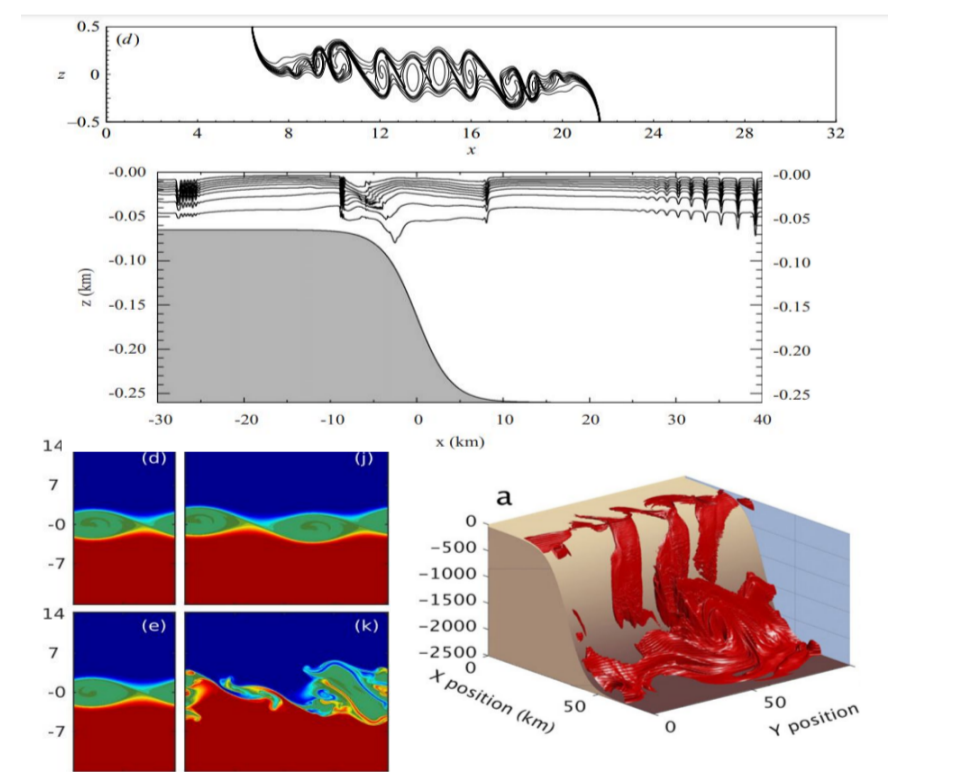
Numerical experiment



Academic Benchmark

	Description	Reference
TC1	Linear acoustic-gravity waves	(Auclair et al. 2021; Harris & Durran 2010)
TC2	Internal tides	(Lamb 2007; Lamb & Kim, 2012)
TC3	Nonlinear internal solitary waves	(Carr et al. 2011; Dunphy et al. 2011)
TC4	Lock exchange (Bq vs Non-Bq)	(Birman et al. 2005)
TC5	Kelvin-Helmholtz instabilities	(Penney et al., 2020)
TC6	Dense shelf overflow	(Yankovsky & Legg 2019)

TC1, TC2, TC3: propagative processes over long distances
TC4, TC5, TC6: local transition to turbulence mechanisms



References

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